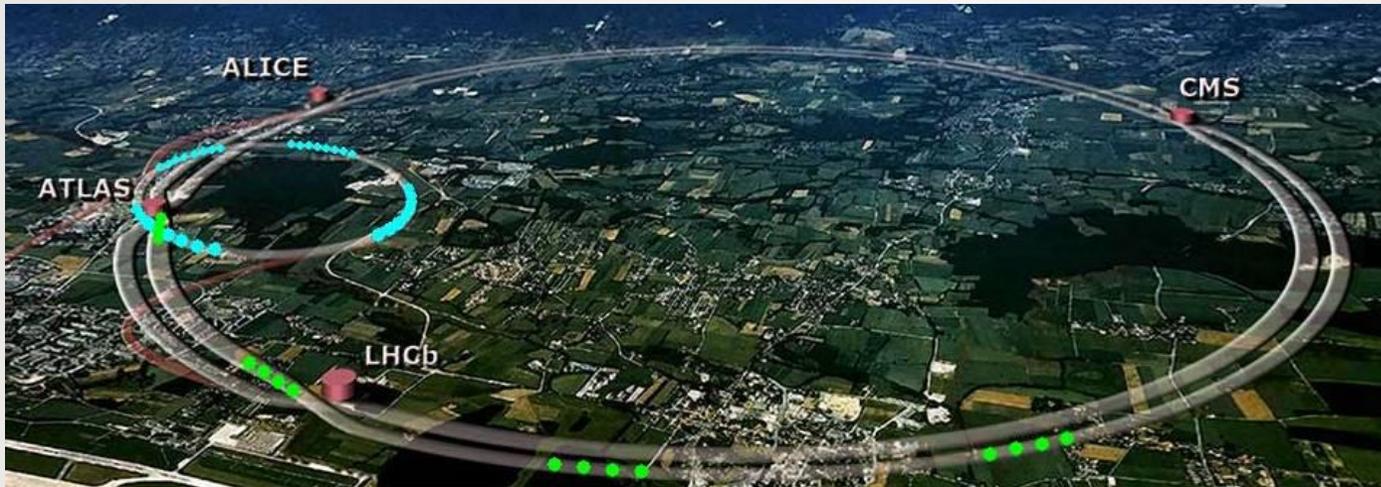


1

# Superconducting Magnet I Conductor



## 2 | Item

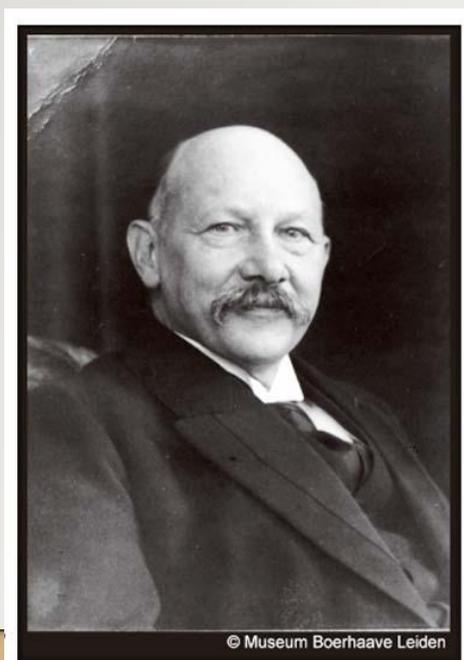
- Superconductivity
- Superconducting Wire and Cables
- Stability and Quench Protection
- Summary

### 3 | Item

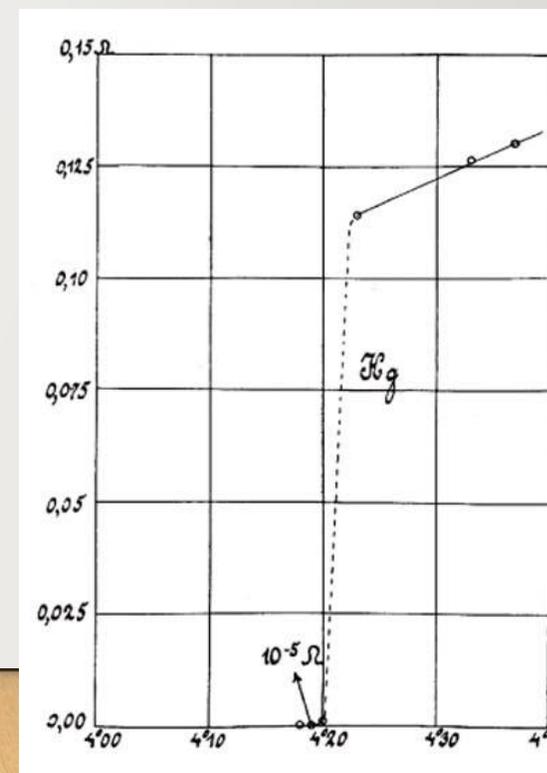
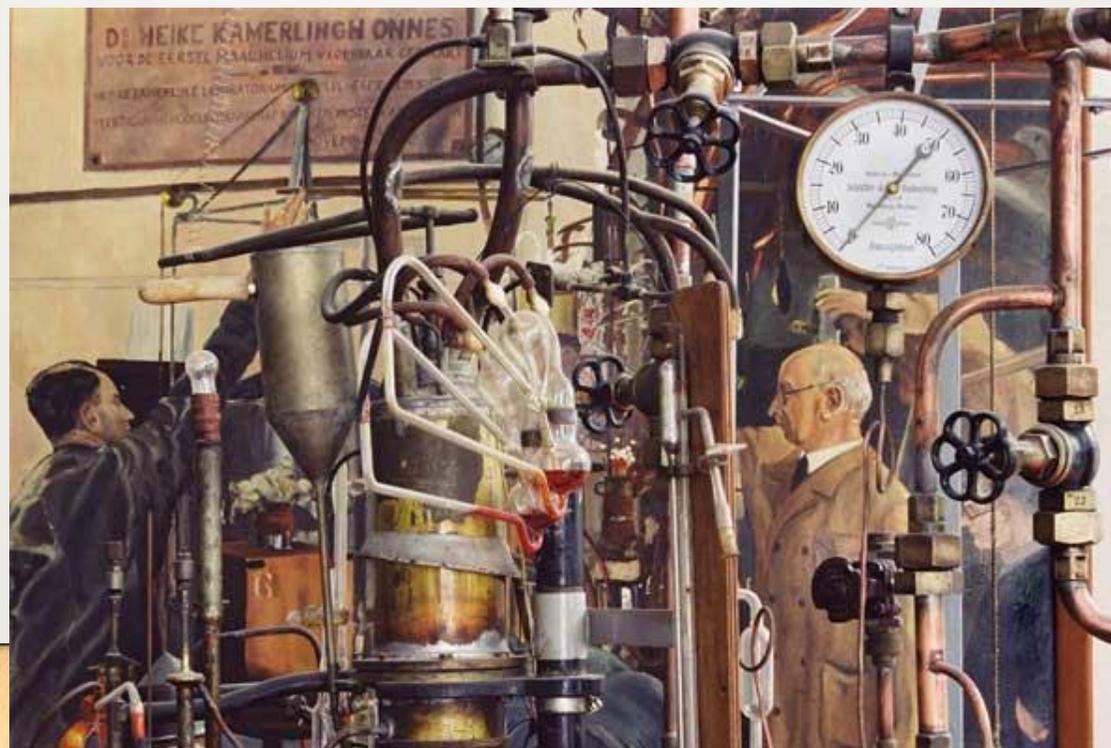
- Superconductivity
- Superconducting Wire and Cables
- Stability and Quench Protection
- Summary

# 4 | 1911 Discovery!

- 1908 Helium liquified
- 1911 April 8th Discovery of Hg Supercon.
- 1913 10T Pb SC magnet proposal  
→ Not realized (Critical Field: 0.08T)



Heike Kamerlingh Onnes (1853-1926)



# 5 | 1933 Meissner Effect

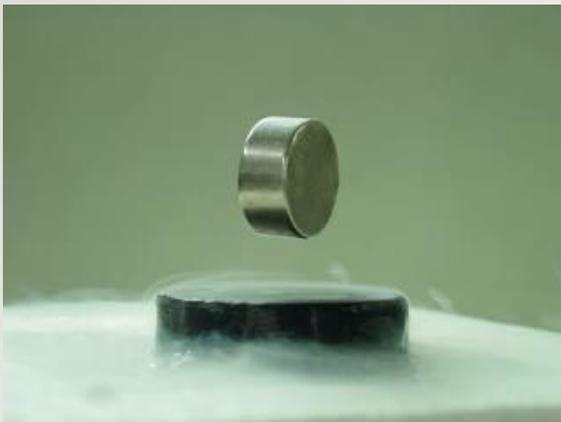
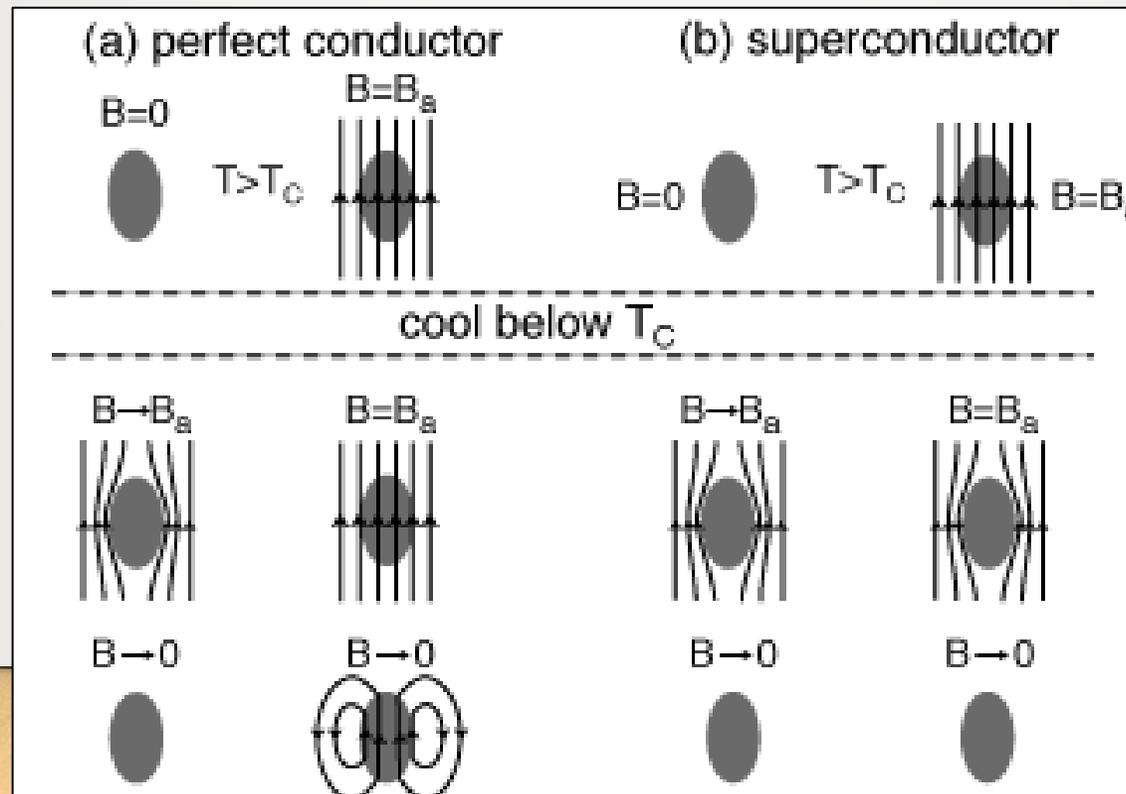
- Found by Fritz Meissner and Robert Ochsenfeld
- Perfect Diamagnetism



Fritz Walther Meissner (1882 –1974)

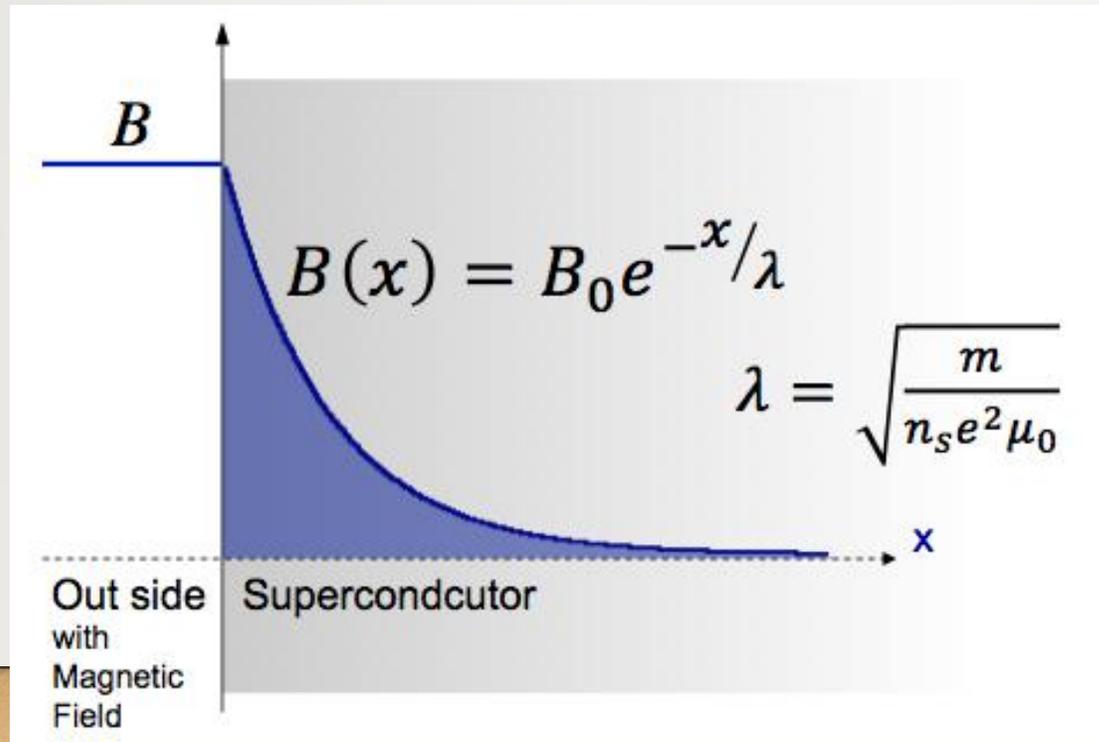


Robert Ochsenfeld (1901–1993)



# 6 | 1935 London Penetration Depth

- Derived by London Brothers
  - Phenomenological Explanation of Meissner effect



Fritz London (1900-1954)

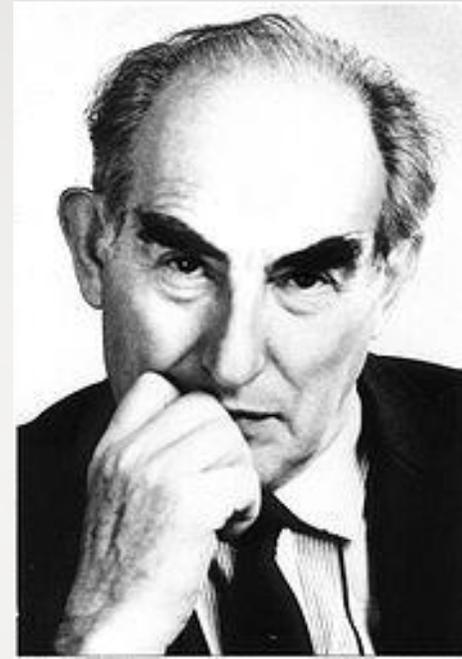


Heinz London (1907-1970)

## 7 | 1950 GL Theory

- Expand London Theory
  - Coherence Length
  - Characteristic length of superconductivity

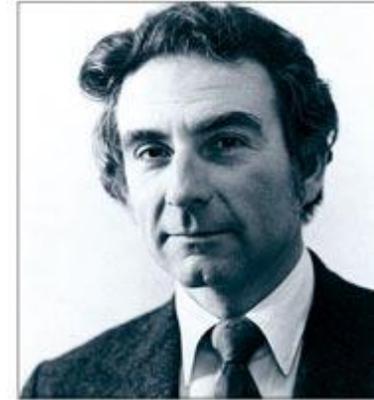
$$\xi = \sqrt{\frac{\hbar}{2m|\alpha|}}, \quad \alpha = \frac{(\mu_0 e \lambda H_c)^2}{m}$$



Vitaly Lazarevich Ginzburg (1916~2009)

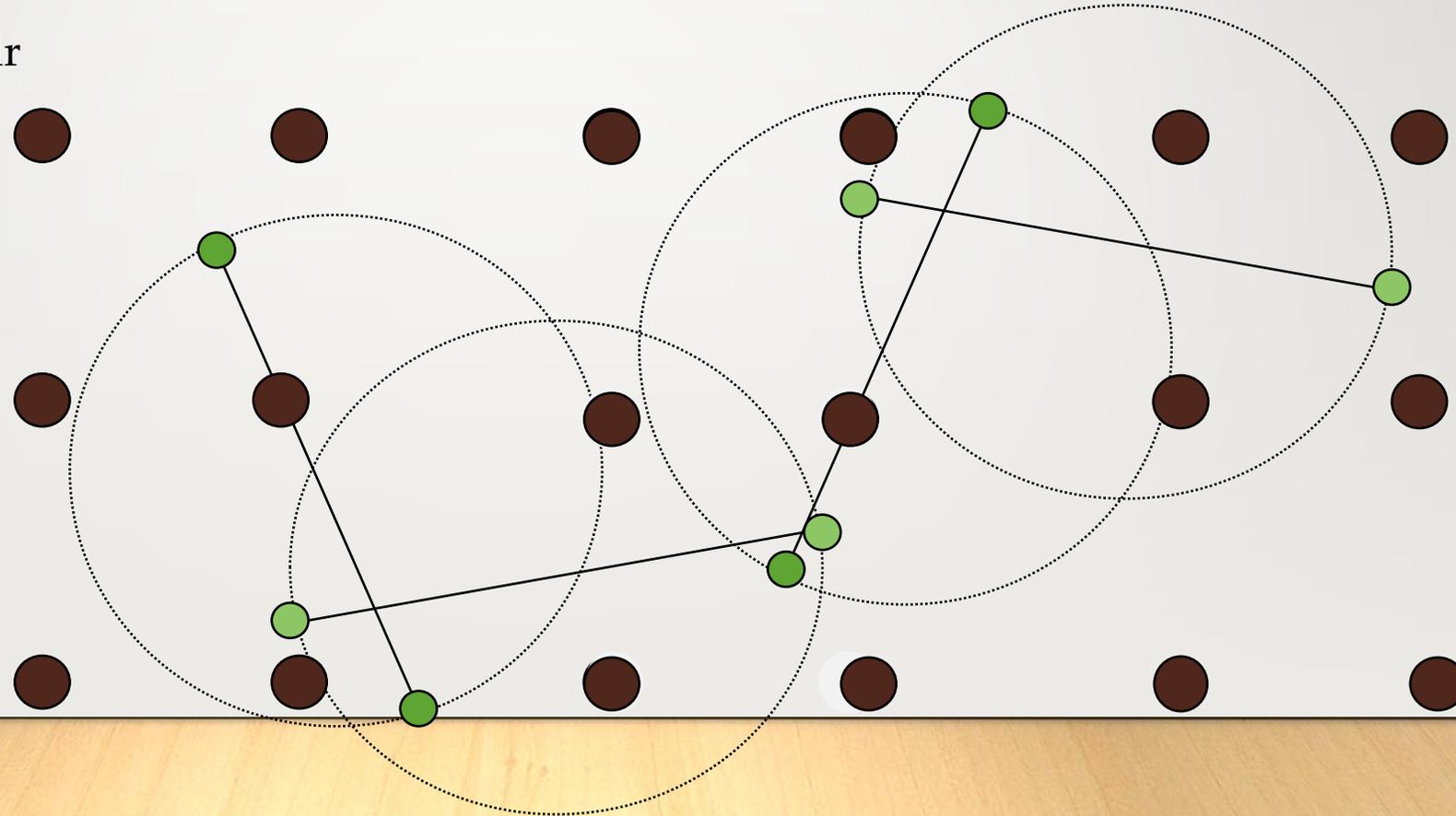


Lev Davidovich Landau (1908~1968)



# 8 | BCS Theory

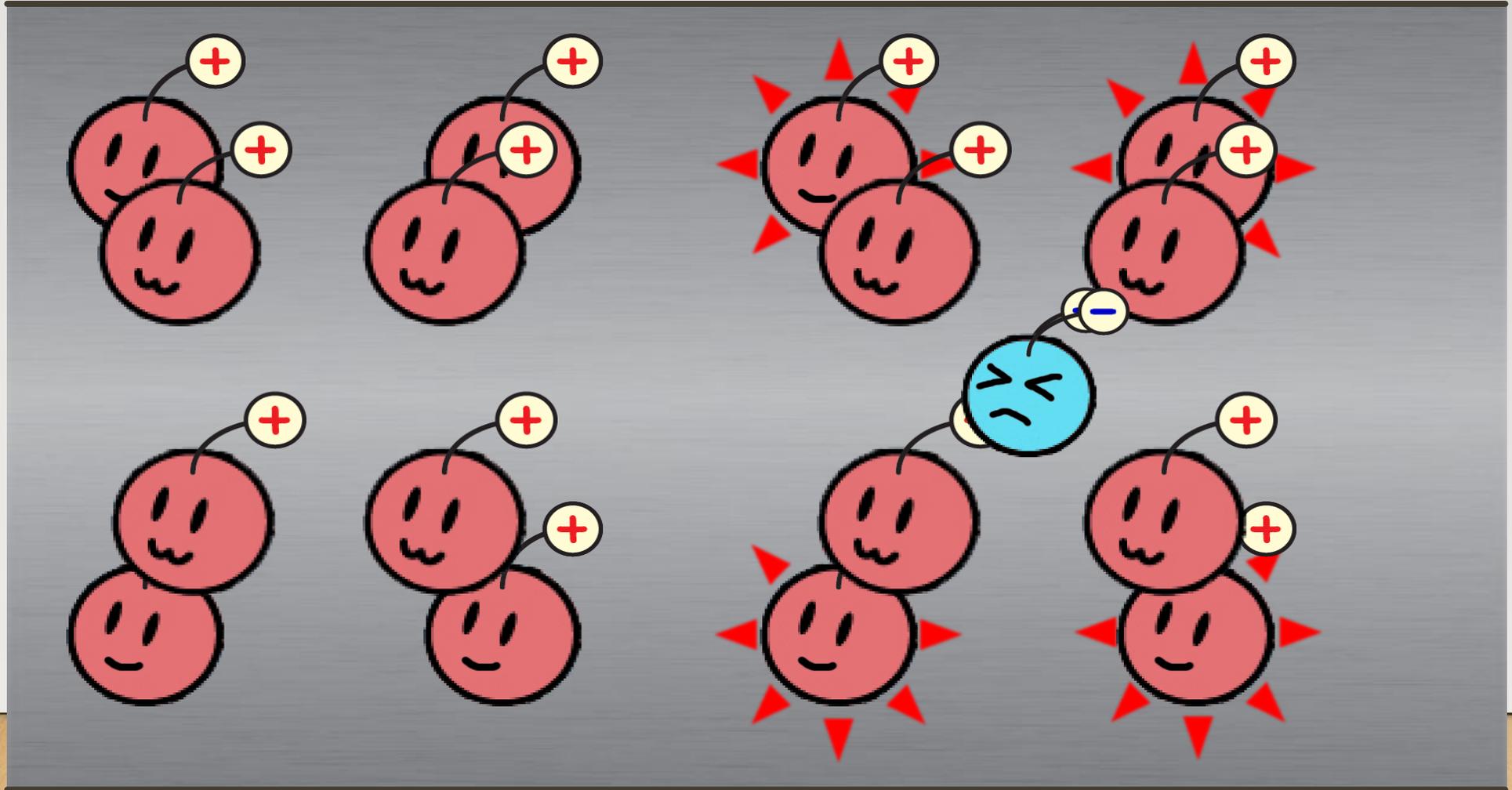
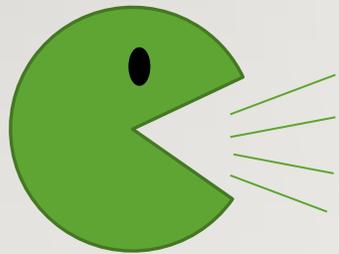
- BCS Theory (1957)
  - Cooper Pair



# 9 | Normal Conductor

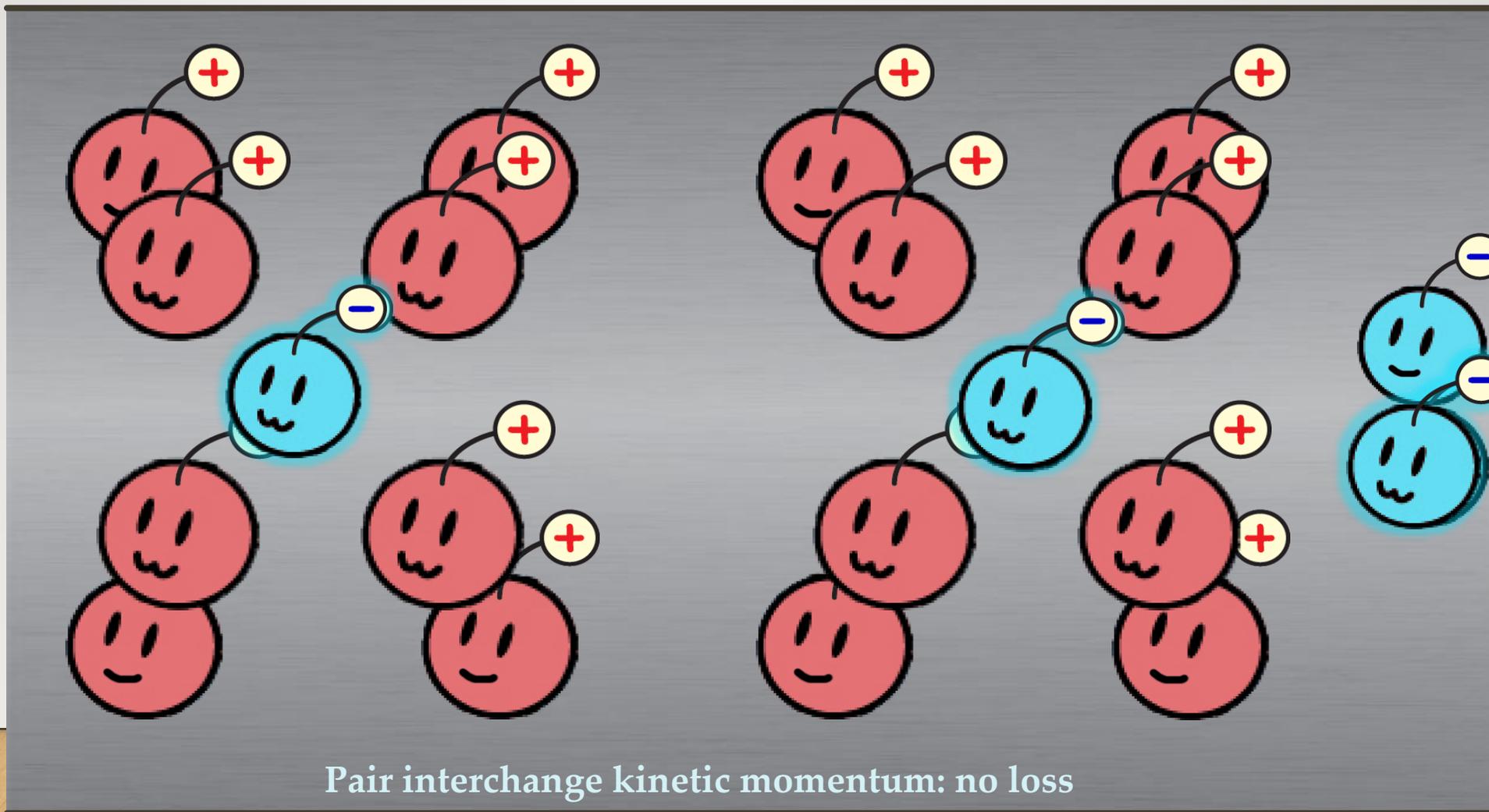
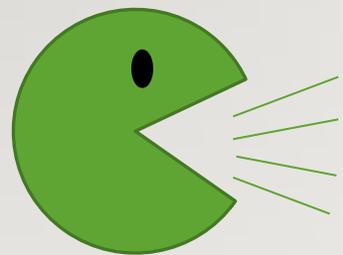
$$W = VI = RI^2$$

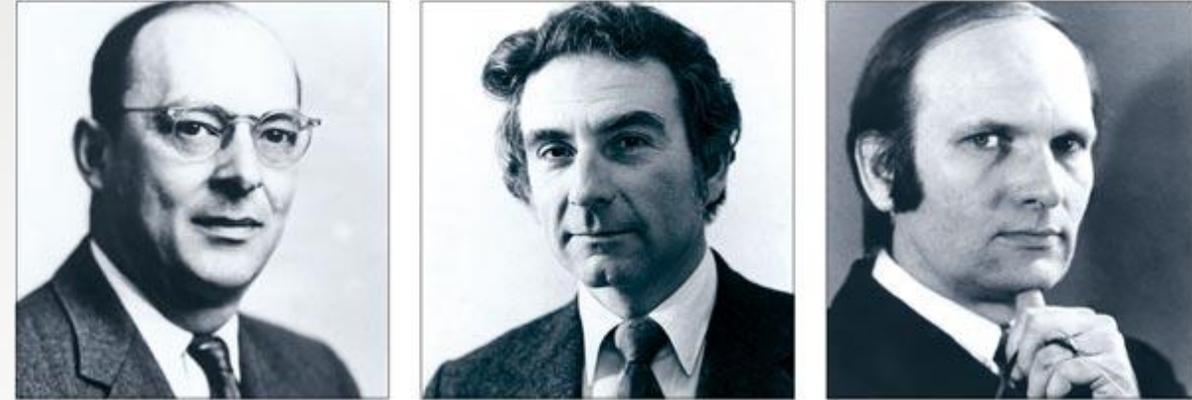
$$V = RI$$



10

# Superconductor

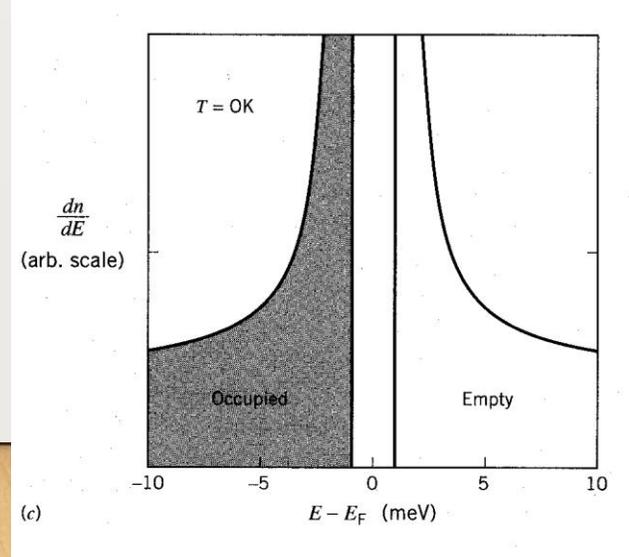
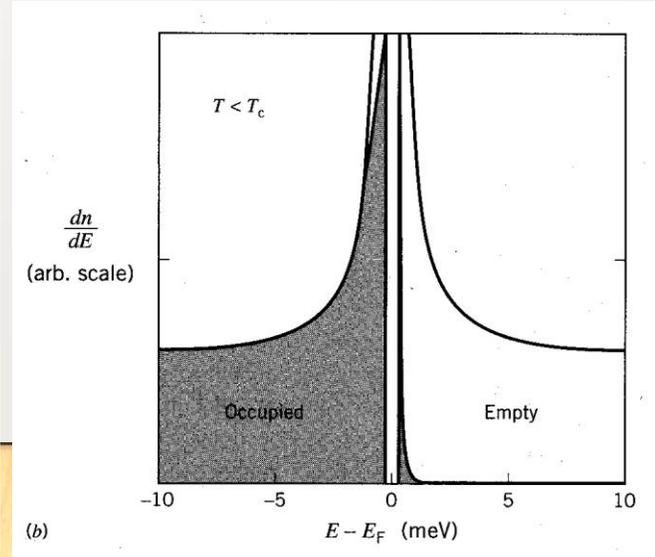
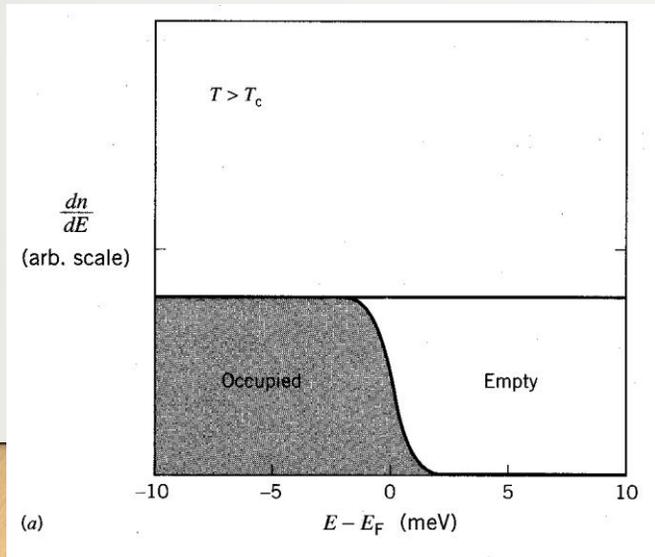
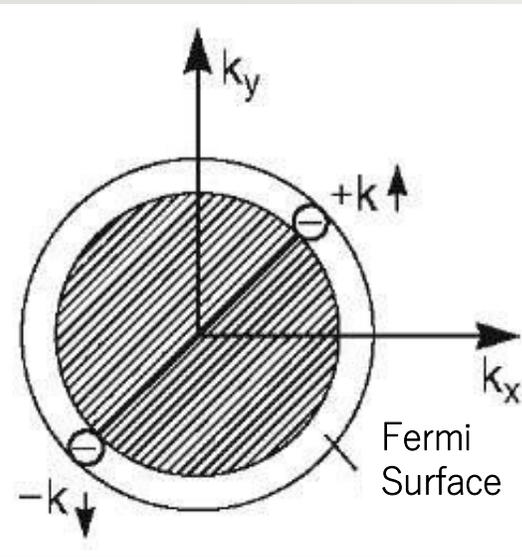




# 11 | BCS Theory

- Based on Cooper pair
- Bose–Einstein condensation
  - All Cooper pairs have the same momentum
  - Energy Gap  $\Delta(T = 0) = 1.764 k_B T_c$ ,
- Second Order Phase Transition
  - Specific Heat Jump at Transition

- Low thermal conductivity
  - No heat transfer by cooper pairs
- Tends to be high resistivity at normal state
  - Large interaction between electron and lattice



12

# Type I Superconductor

Type I superconductor

$$\xi \gg \lambda$$

Positive surface energy



Decrease boundary



Expel flux

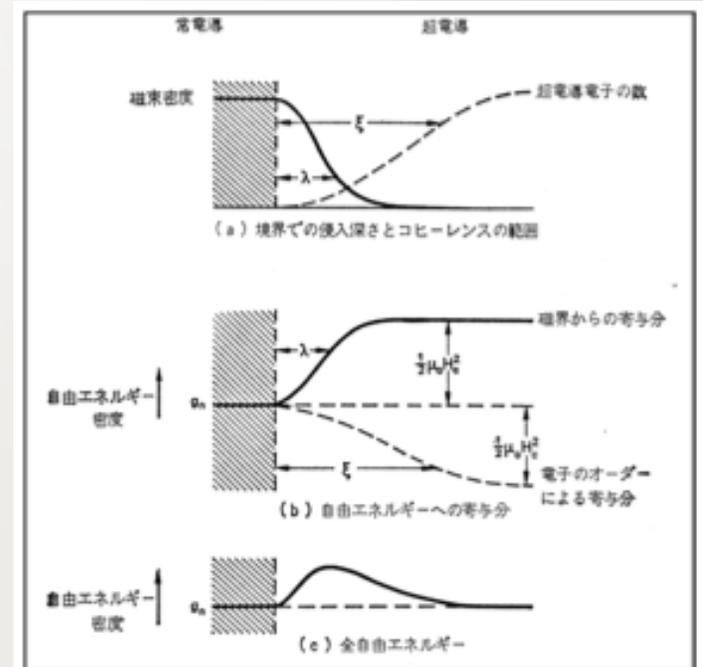
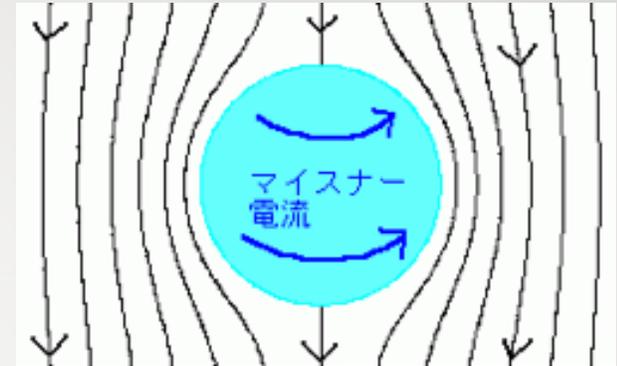
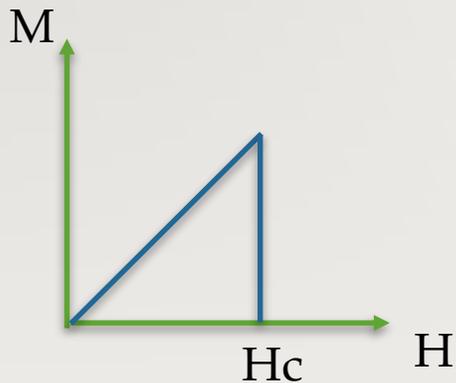


Fig. 2-3 第1種超伝導体の界面[18]

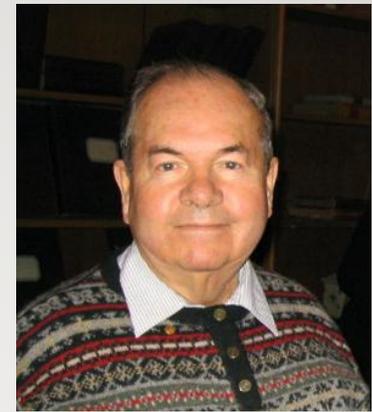
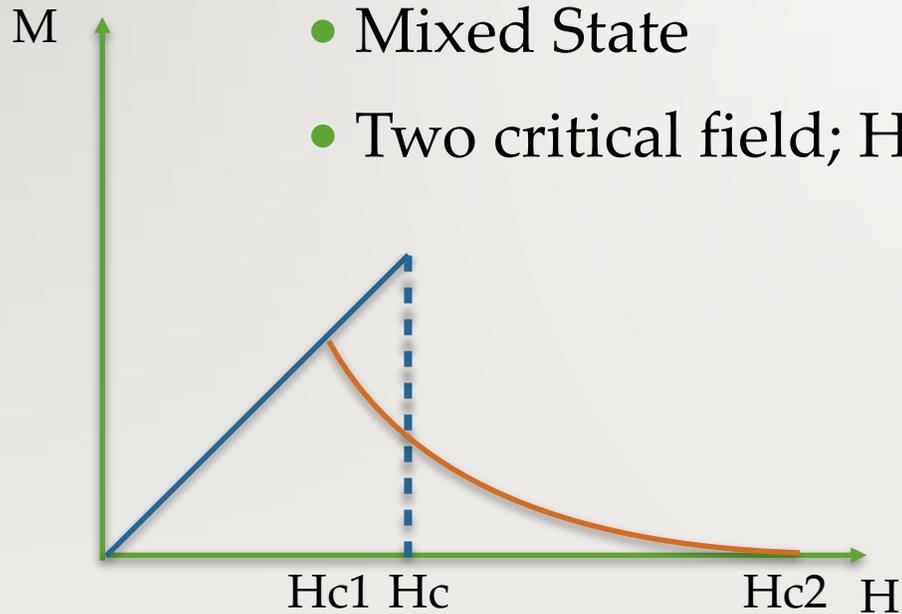
13

# 1957 Type II Superconductor

- Derived by Abrikosov from GL theory

$$\sqrt{2} \xi < \lambda$$

- Boundary Surface Energy: Negative
- Mixed State
- Two critical field;  $H_{c1}$ ,  $H_{c2}$



Alexei Alexeyevich Abrikosov

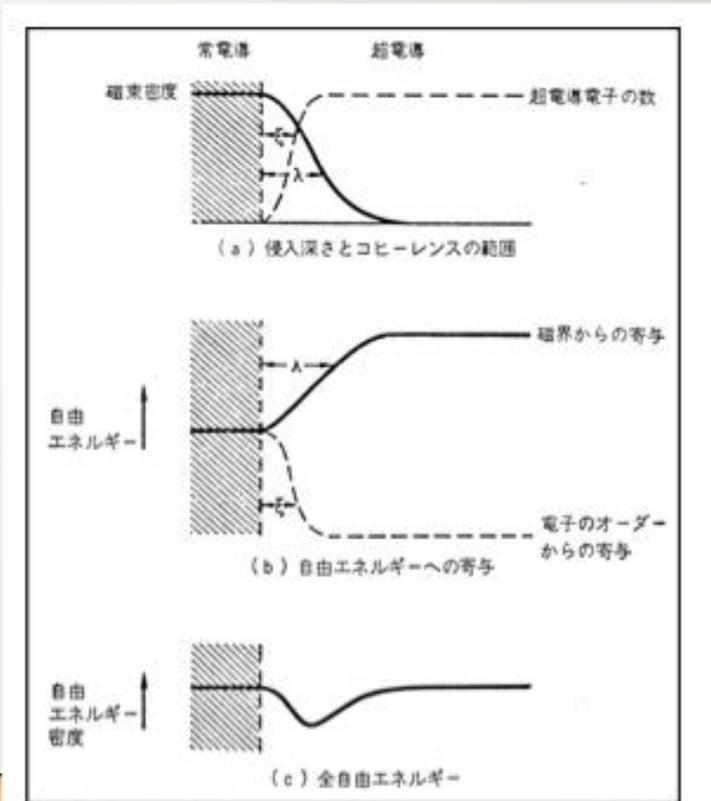
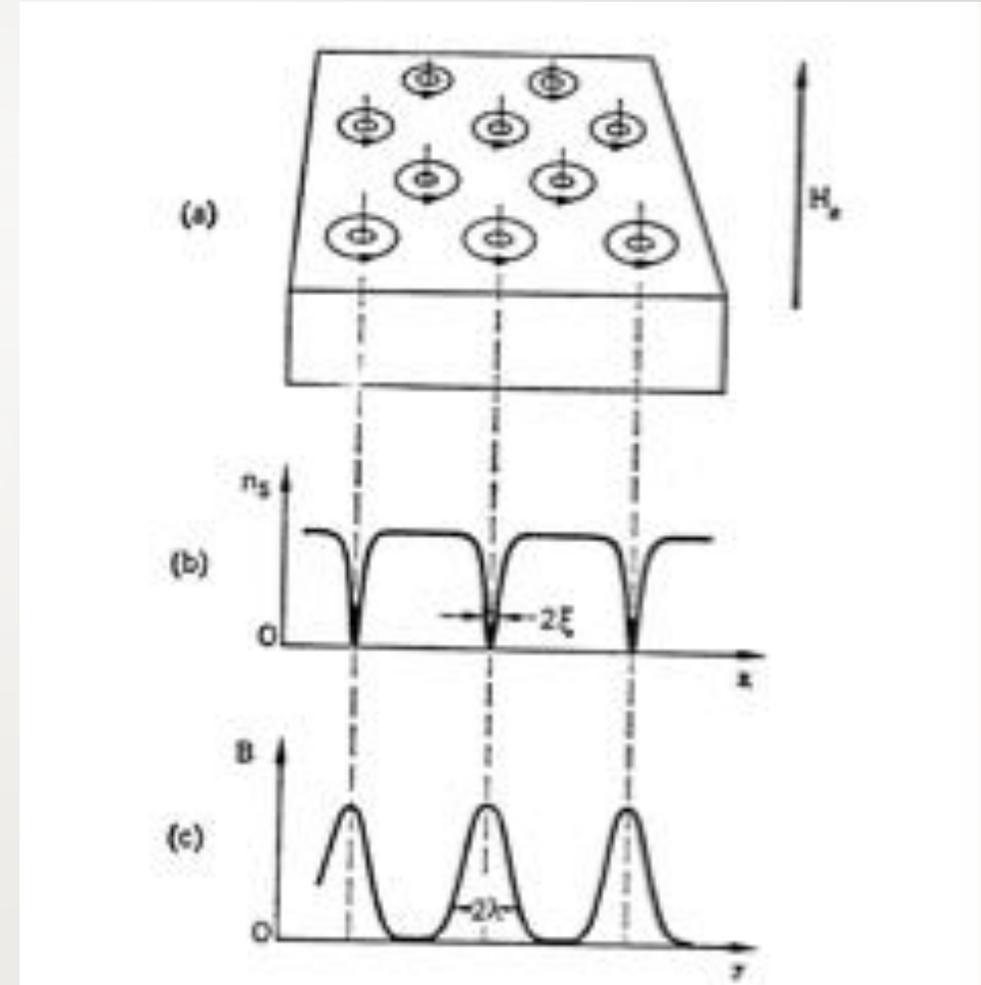
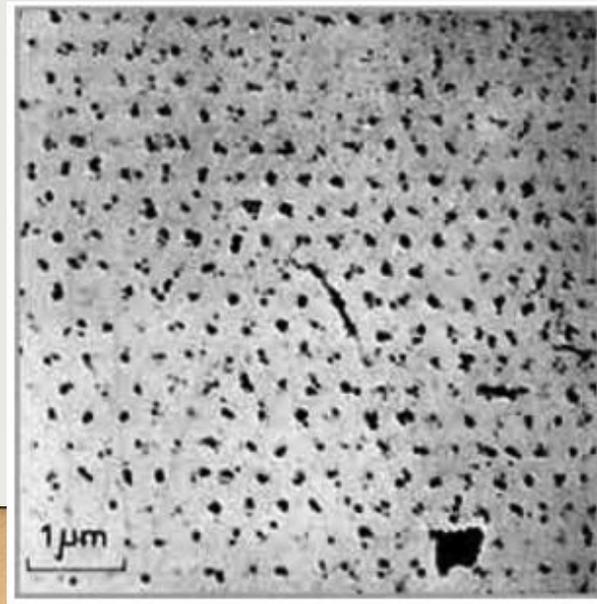


Fig. 2-4 第2種超伝導体の界面[18]

# 14 | Mixed State

- Negative Surface Energy
- Maximize NS boundary
- minimum flux
  - Cooper Pair: Vortex



15

## Type II Superconductor

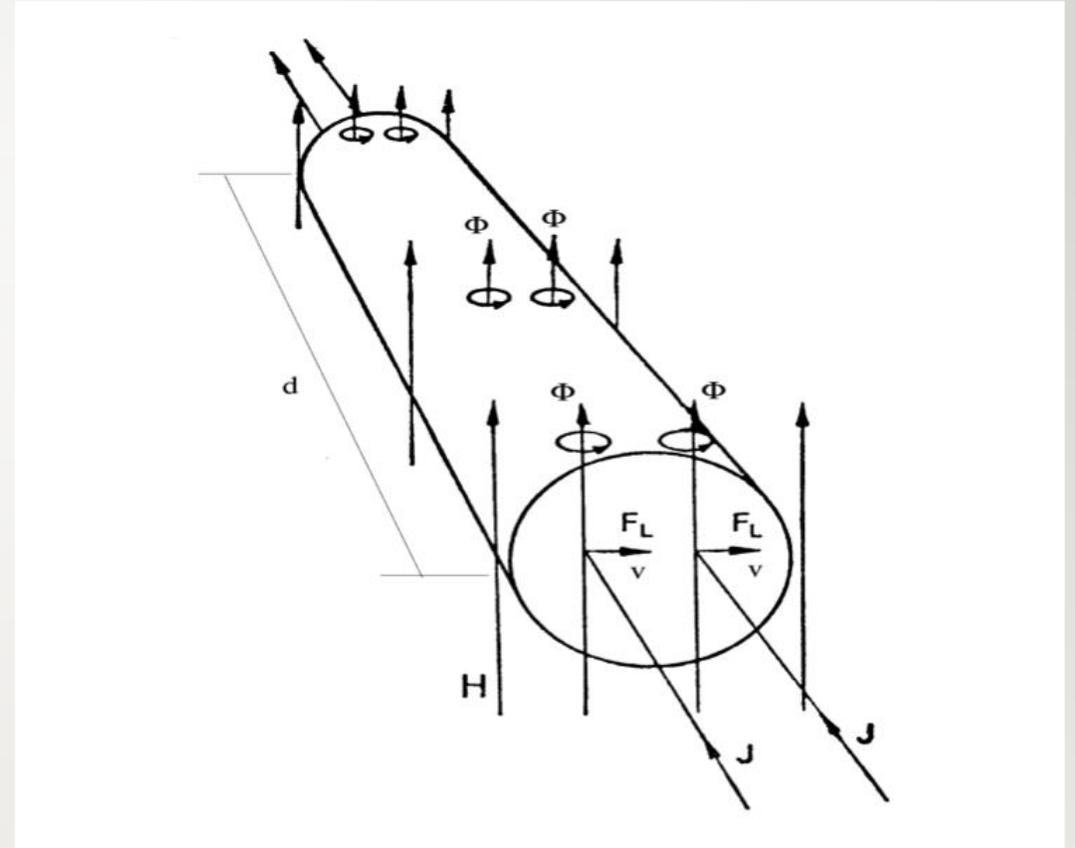
- Coherence Length by Pippard

$$\frac{1}{\xi} \approx \frac{1}{\xi_0} + \frac{1}{L}$$

- where  $\xi_0 \approx \frac{\hbar v_F}{\pi \Delta(0)} \approx 0.18 \frac{\hbar v_F}{k_B T_c}$  is coherence length of ideal material,  
 $L$  is electron mean free path
- Dirtier Material: shorter coherence length: type II conductor

# 16 | Flux Flow

- Field  $\times$  Current = Lorentz Force
- Flux may move: Flux Flow
- Flux Flow = Electric Field = Loss

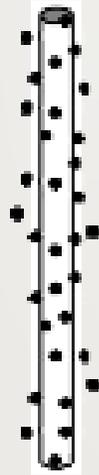


$$V = n\Phi vd$$

Voltage along current direction

# 17 | Pinning and Critical Current

- Defect: easier to be normal state
  - Create Pinning Potential
- Type of Defects



Point Defects



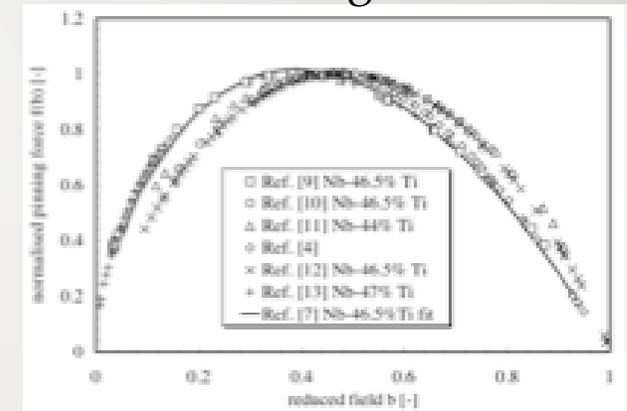
Dislocation



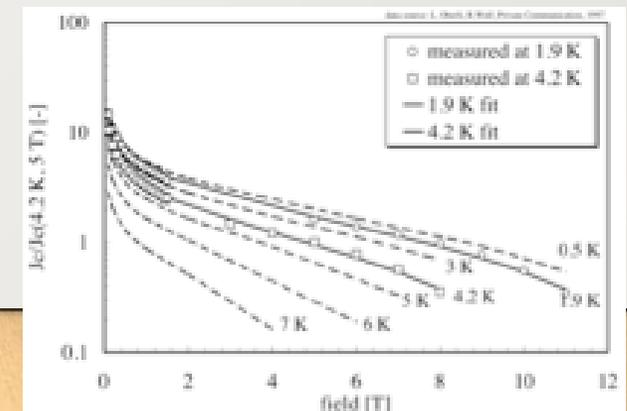
Grain Boundary

Pinning Potential

Pinning Force

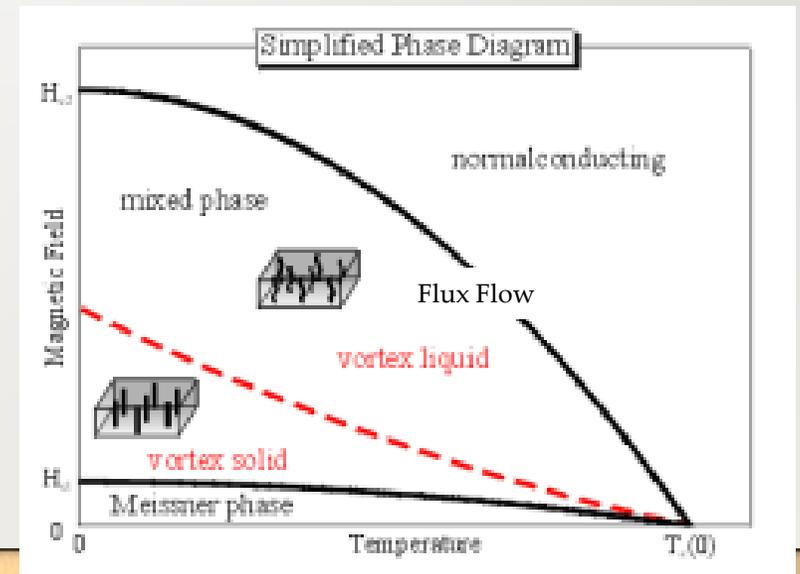
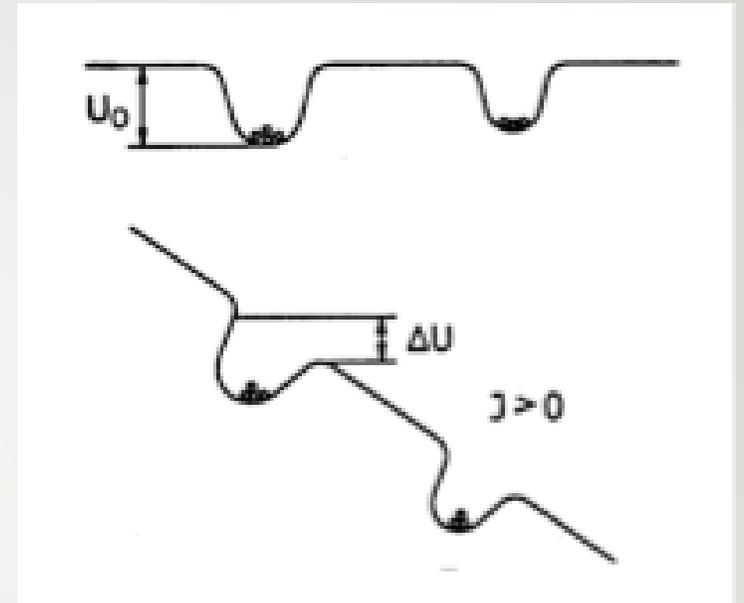


Critical Current Density



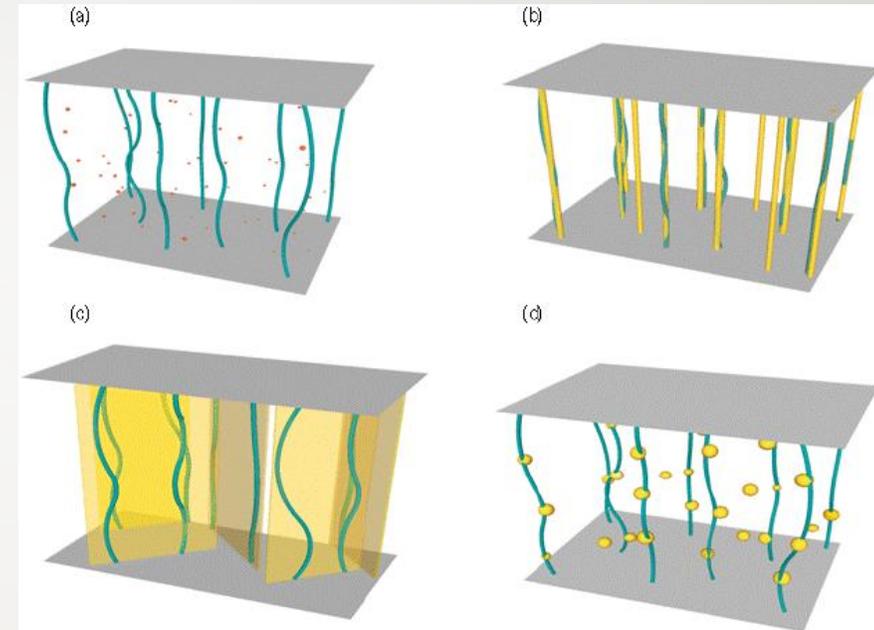
# 18 Flux Creep, Melting & Flow

- Flux pinning potential:  $U_0$ 
  - Lorentz force potential difference:  $\Delta U$
- Flux Creep ( $U_0 > \Delta U$ , T: Low)
  - Thermal excitation of flux
  - Flux move: Logarithmic decay
- Flux Melting ( $U_0 > \Delta U$ , T: High)
  - Higher Temperature
  - Larger excitation of flux
  - Flux move: Continuous
- Flux Flow ( $U_0 < \Delta U$ )
  - Flux move: Continuous



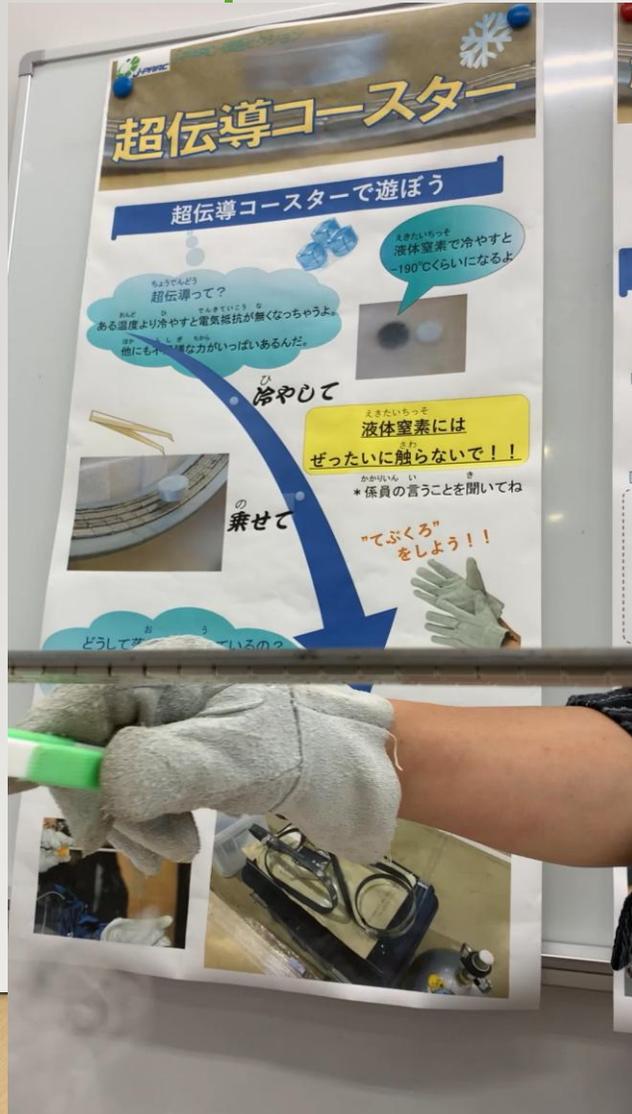
# 19 | Artificial Pinning

- To Create Deeper Pinning Potential
- Artificially insert normal conducting nano-structure
  - Nano Rod
  - Nano Particle



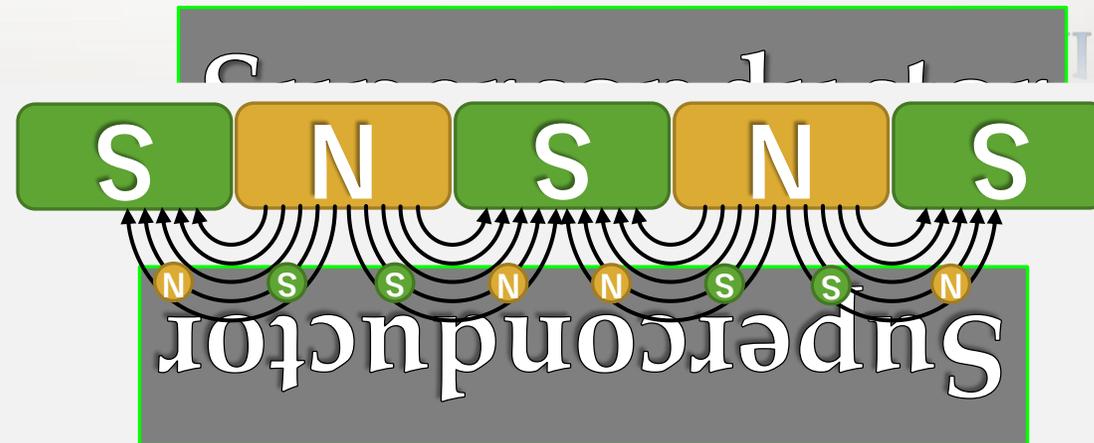
Sketch of the dimensionality of several artificial pinning centers (APCs) with examples of typical defects for each case. (a) OD-PC, oxygen vacancies, element substitutions, point defects (b) 1D-PC such as dislocations and columnar defects, (c) 2D-PC, like grain boundaries, twin boundaries or planar defects and (d) 3D-PC, precipitates, secondary phases or local strain.

# 20 | Meissner Effect and Pinning Effect



## Stick by Pinning Effect

Enforce flux pinning by pushing superconductor to magnet rail

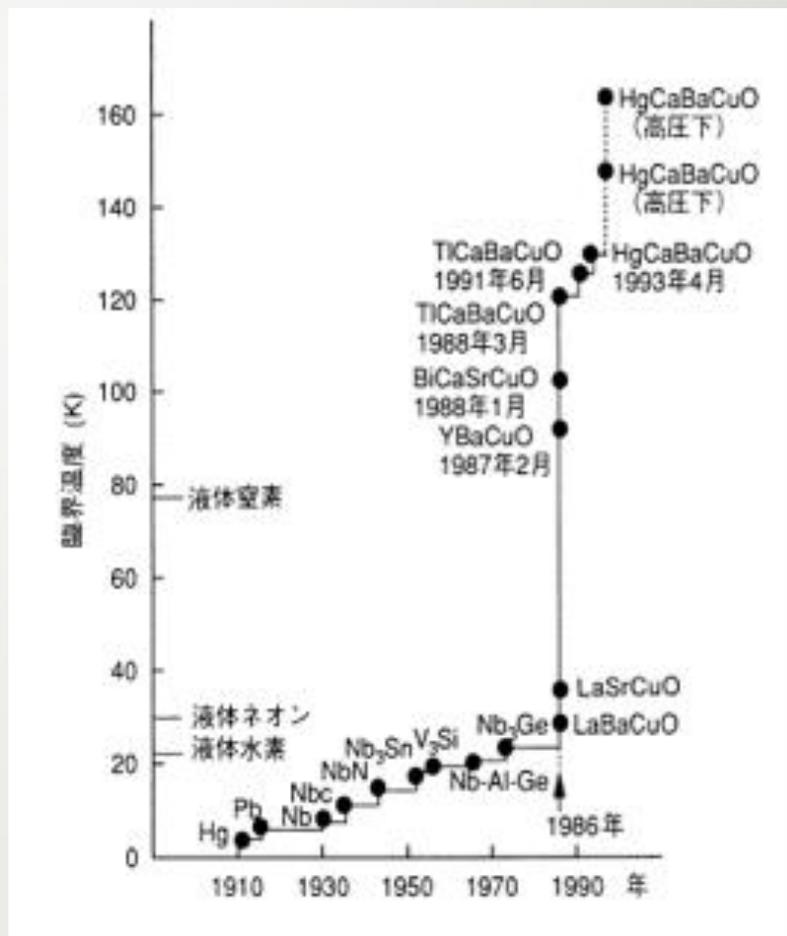


Cross Section

Float by Meissner Effect

# Superconductor

- Type I
  - Hg 4.153K
  - Pb 7.163K
- Type II
  - Nb 9.46K
  - Nb<sub>3</sub>Sn 18.1K
  - MgB<sub>2</sub> 39K
  - YBa<sub>2</sub>Cu<sub>3</sub>O<sub>7</sub> 92K



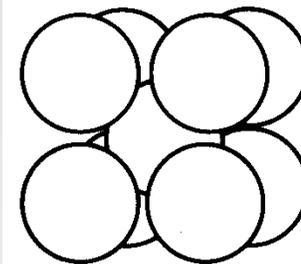
## 22 | Superconductor Application

- for Superconducting RF
  - Meissner state, high  $H_{c1}$
  - Pure, Clean, no Defects
- for Superconducting Magnets
  - Mixed State, high  $H_{c2}$
  - Strong Pinning Force, high  $J_c$
  - Dirty and lot of Defects

- Superconductivity
- Superconducting Wire and Cables
- Stability and Quench Protection
- Summary

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# Superconductors

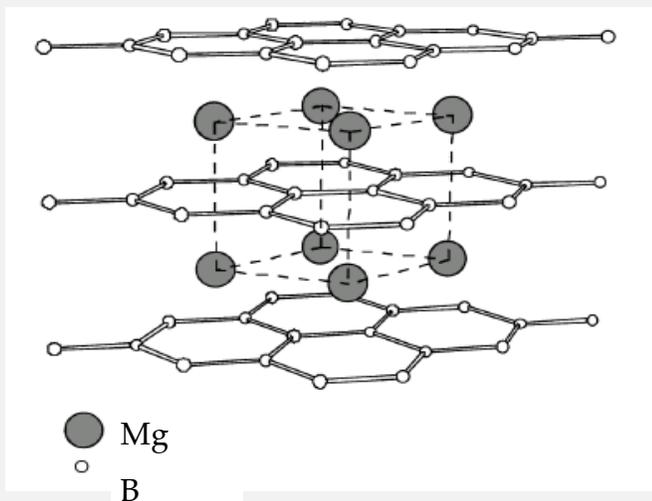


Alloy

Tc: 9K

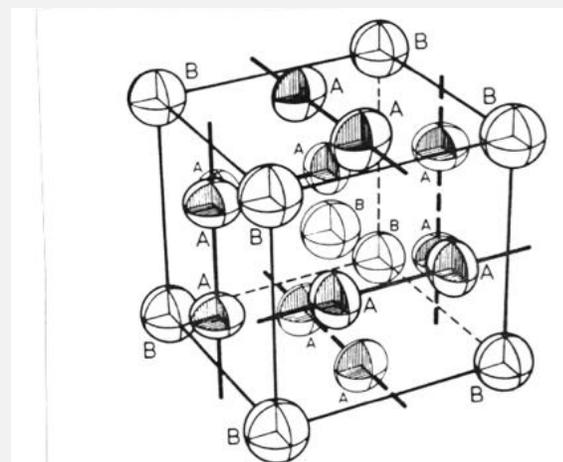
Nb-Ti

ductile



Tc:39 K

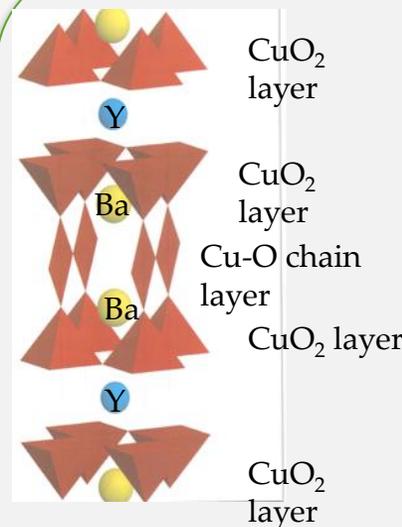
MgB<sub>2</sub>



Tc:18 K

Nb<sub>3</sub>Sn

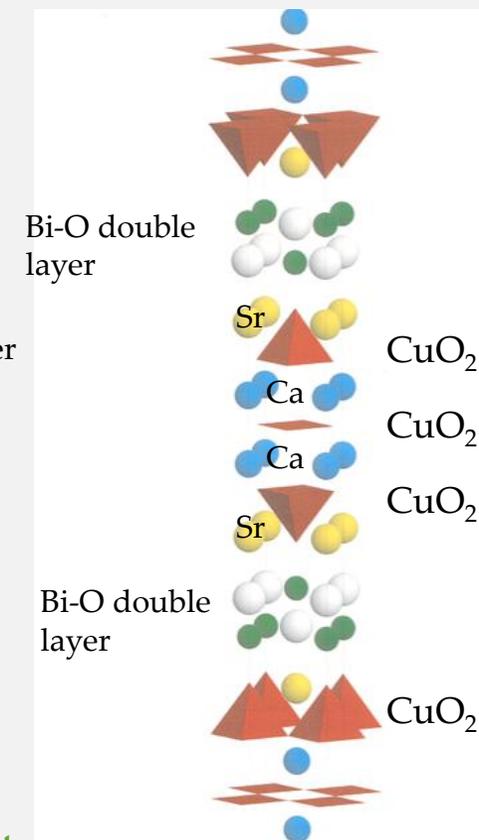
Metallic Compound



Tc: 92-95 K

YBCO

Cuprate-perovskite  
Ceramic



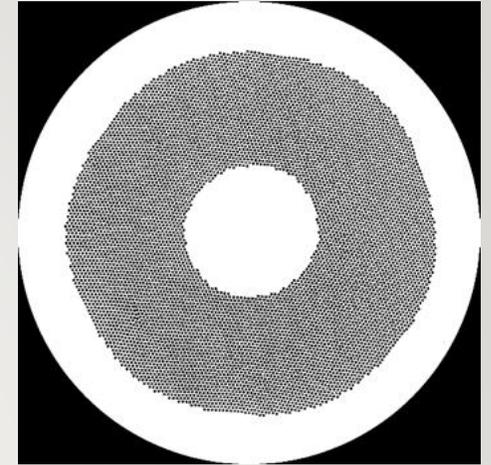
Tc:110 K

BISCO

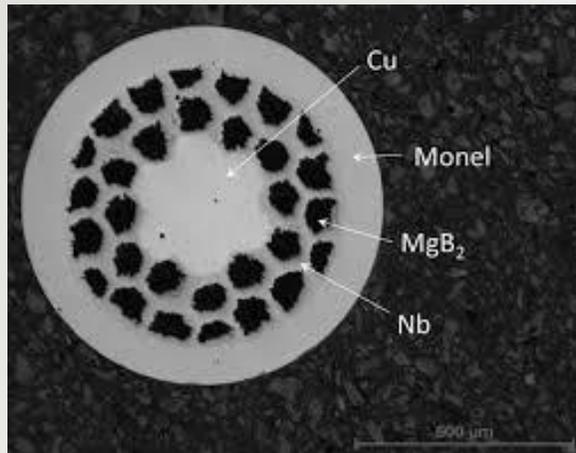
25

# Actual Superconductors

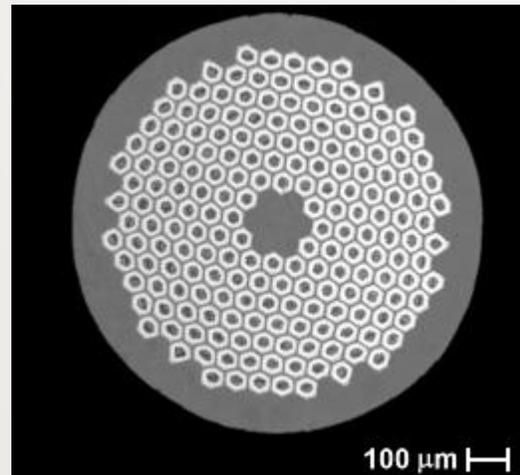
Nb-Ti



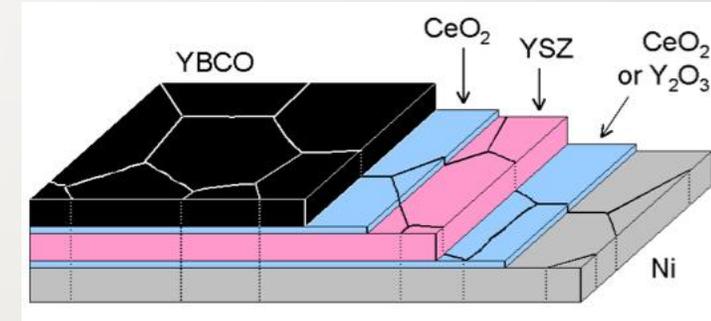
MgB<sub>2</sub>



Nb<sub>3</sub>Sn



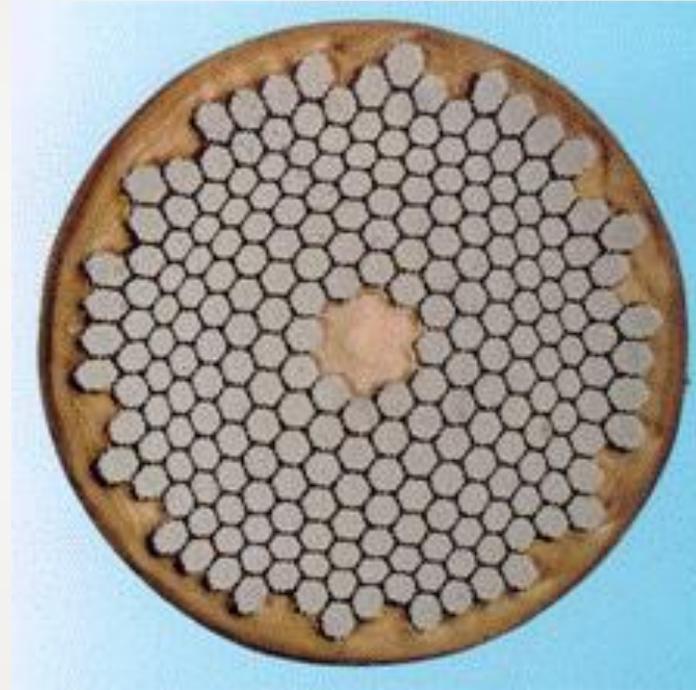
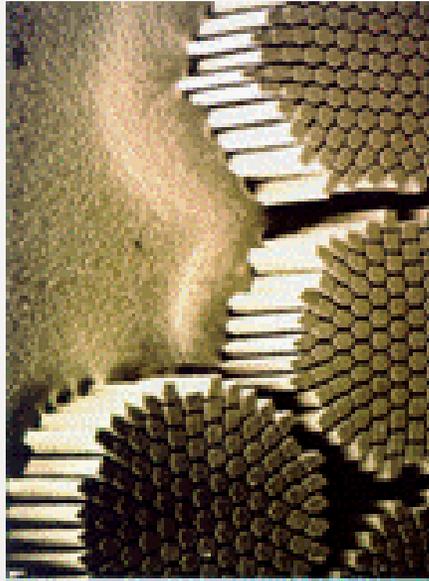
Coated conductor. (IBAD, RABiTS or ISD)



Bi-2223



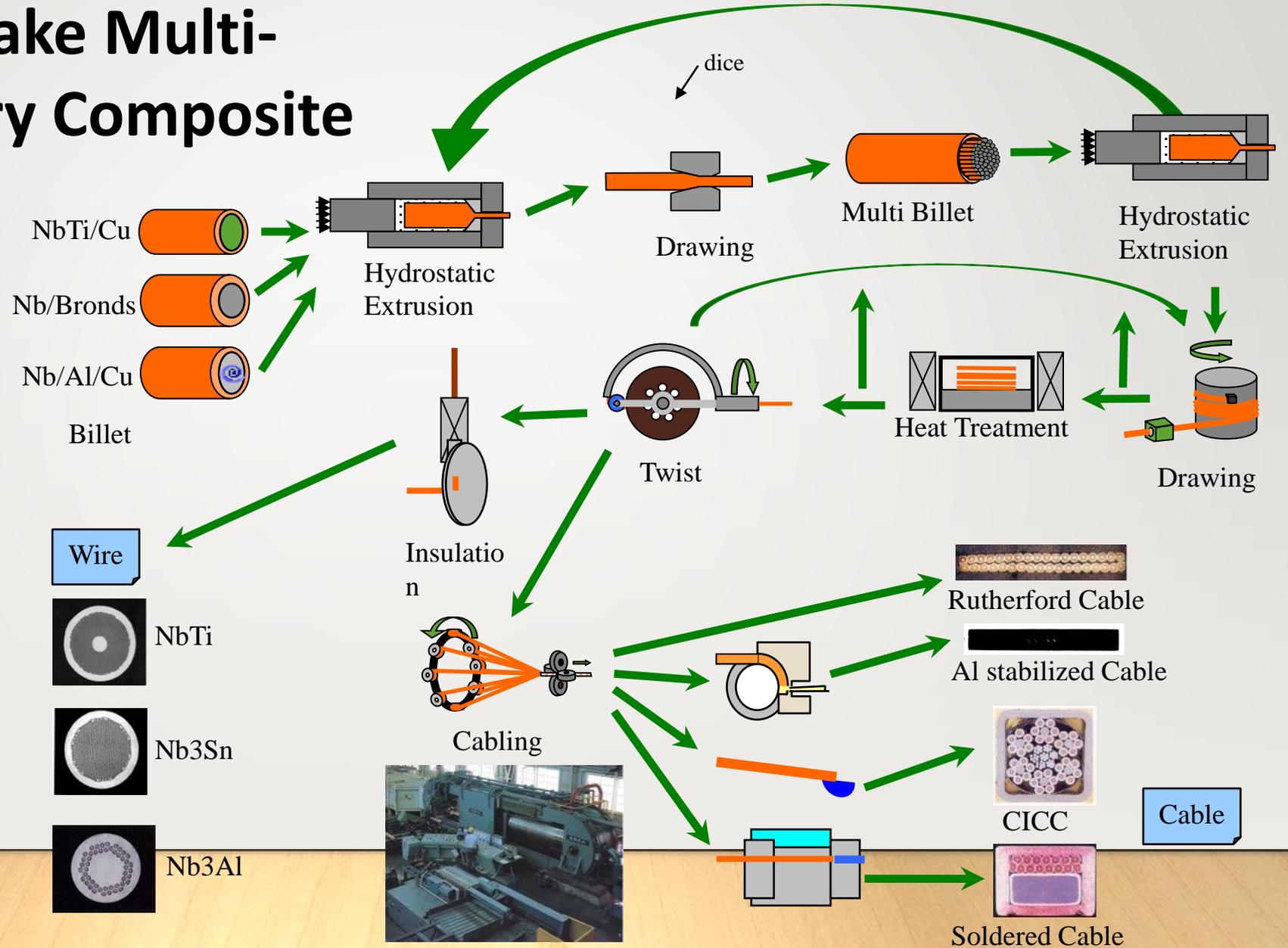
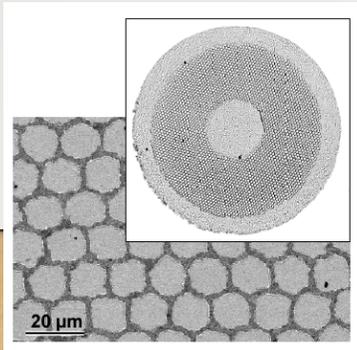
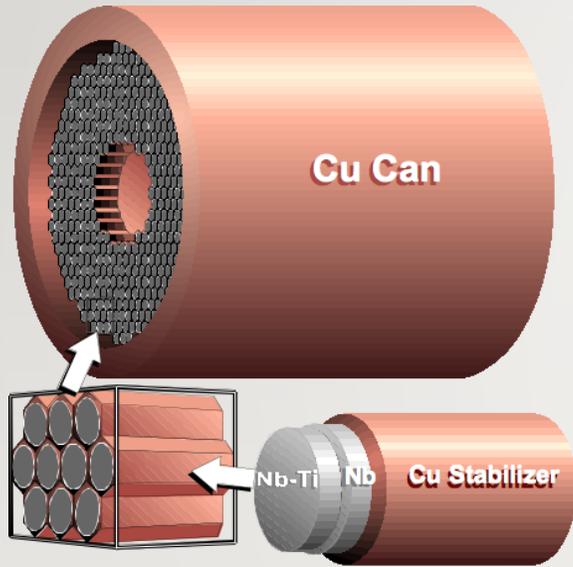
# NbTi Multi-filamentary Composite Wire (Strand)



- Fine multi-filament embedded in copper matrix
  - Filament Diameter : 100~1  $\mu\text{m}$ 
    - Depending on purpose

27

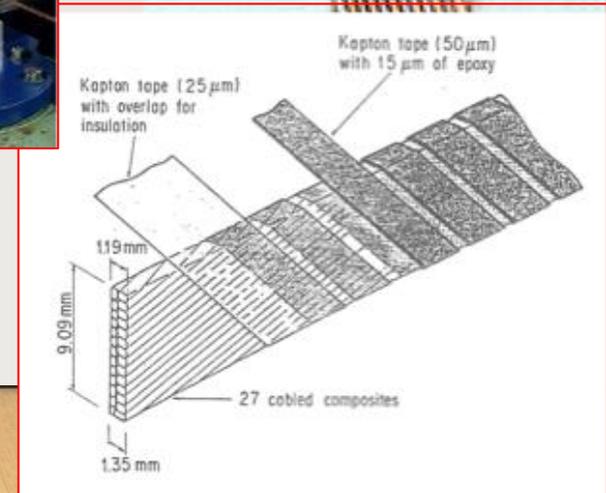
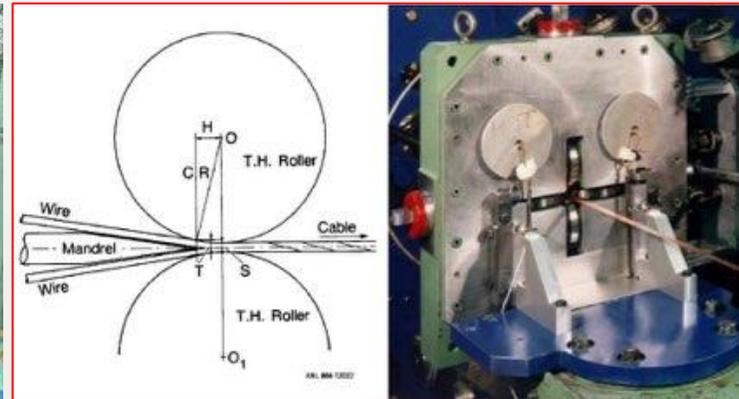
# How to Make Multi-filamentary Composite



28

# Rutherford Cable

- 10~60 Strand twisted together into flat squad cable
- Eventual Keystone angle



# Why such a complicated structure?

- **Improve Critical Current Density**
  - **Pinning Center**
- **Optimize Magnetization Properties**
  - **Minimize Diamagnetic Magnetization of Superconductor**
    - Minimize Errors to Magnetic Field (Class 2)
    - Minimize AC-Loss by Hysteric Loss
- **Stabilization**
  - **Minimize Disturbance**
    - Magneto-Thermal Instability = Flux Jump
  - **Stabilize Against Disturbances**
    - Improve electric conductivity at normal state: less ohmic heating
    - Improve thermal conductivity: better cooling
- **Quench Protection**
  - **Guarantee safe shut down of magnet**
    - Improve electric conductivity at normal state: less ohmic heating
    - Increase heat capacity: Less temperature increase
    - Improve thermal conductivity: better heat distribution

# Critical Current Density

- Critical Current Density



- Pinning Force

$$F_{pc} = n\Phi J_c = BJ_c$$

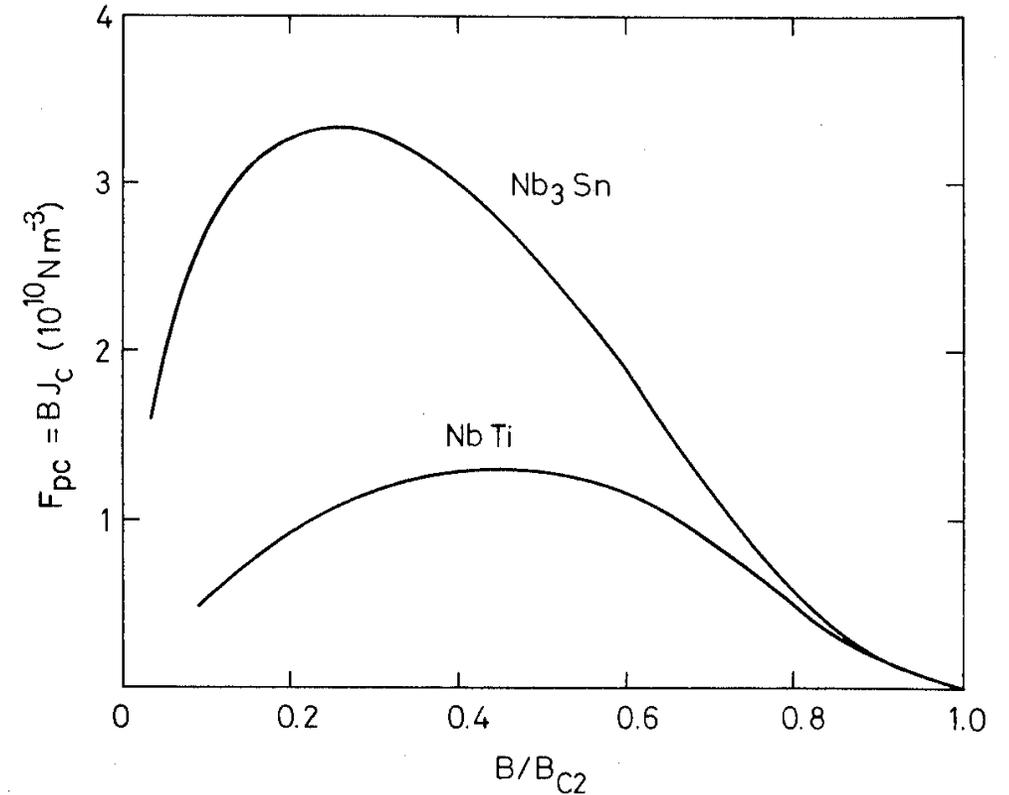


Fig. 12.3. Pinning force per unit volume, as a function of reduced field, for niobium titanium and niobium tin.

# Introducing Pinning Centers

Drawing Process

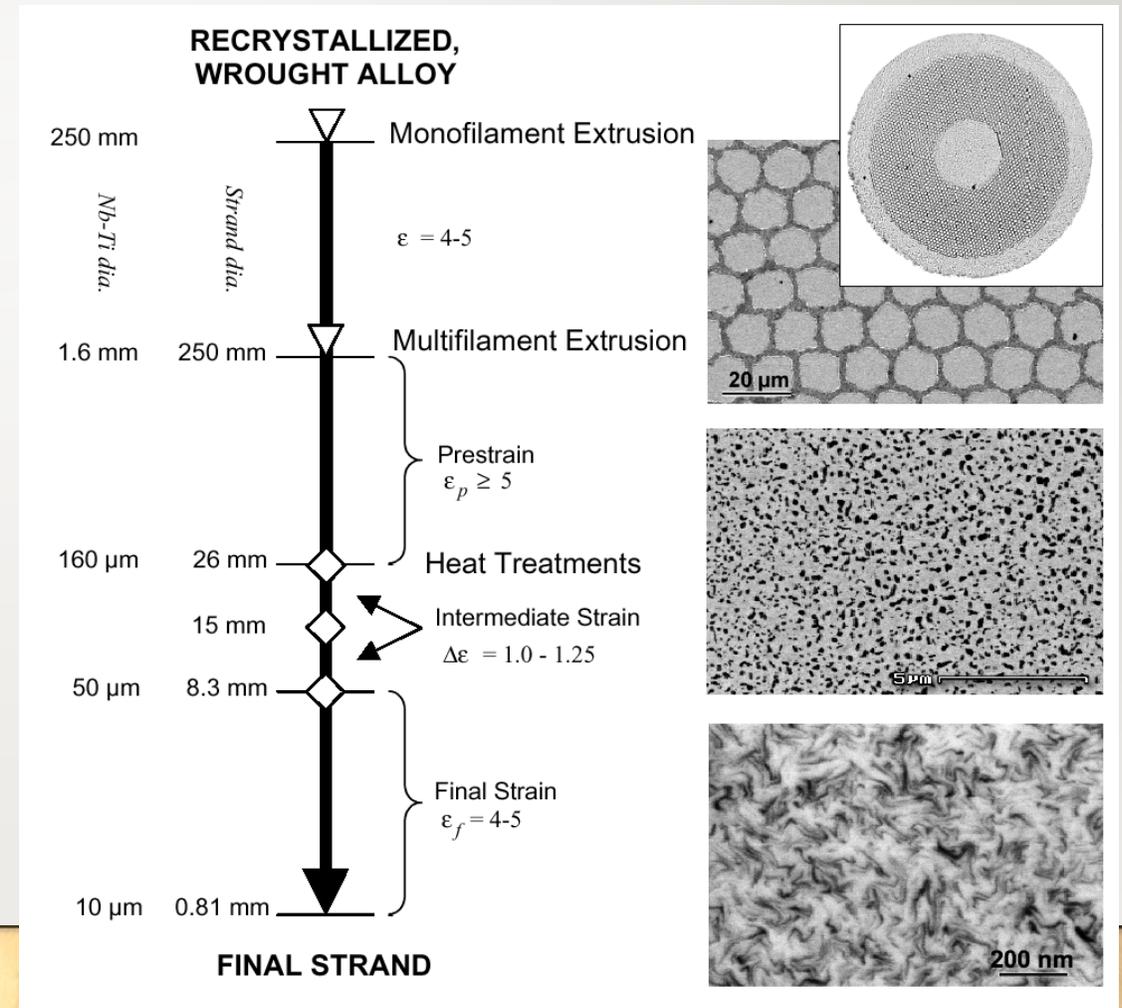
Draw



Heat Treatment



Lattice Defects

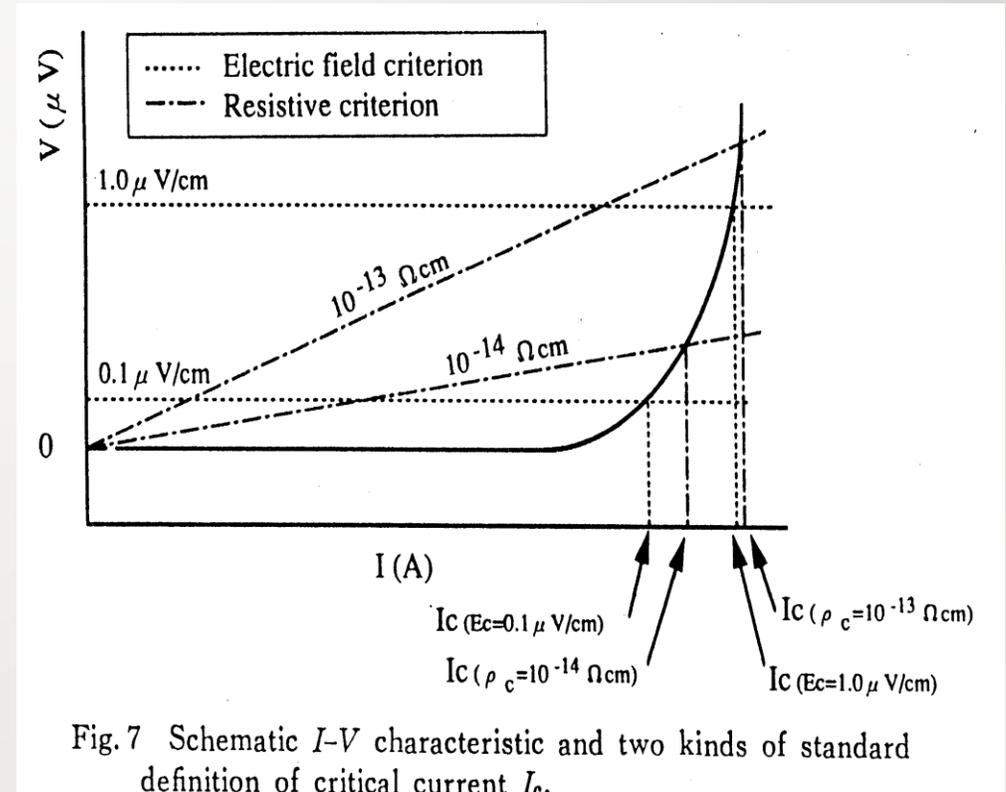


# Introducing Pinning Centers?

- Conductor electric characteristic

$$V \propto I^n$$

- n-value
  - Bad n-value
  - = bad critical current
  - Filament Damage?
- Definition of critical current
  - Multiple Definition





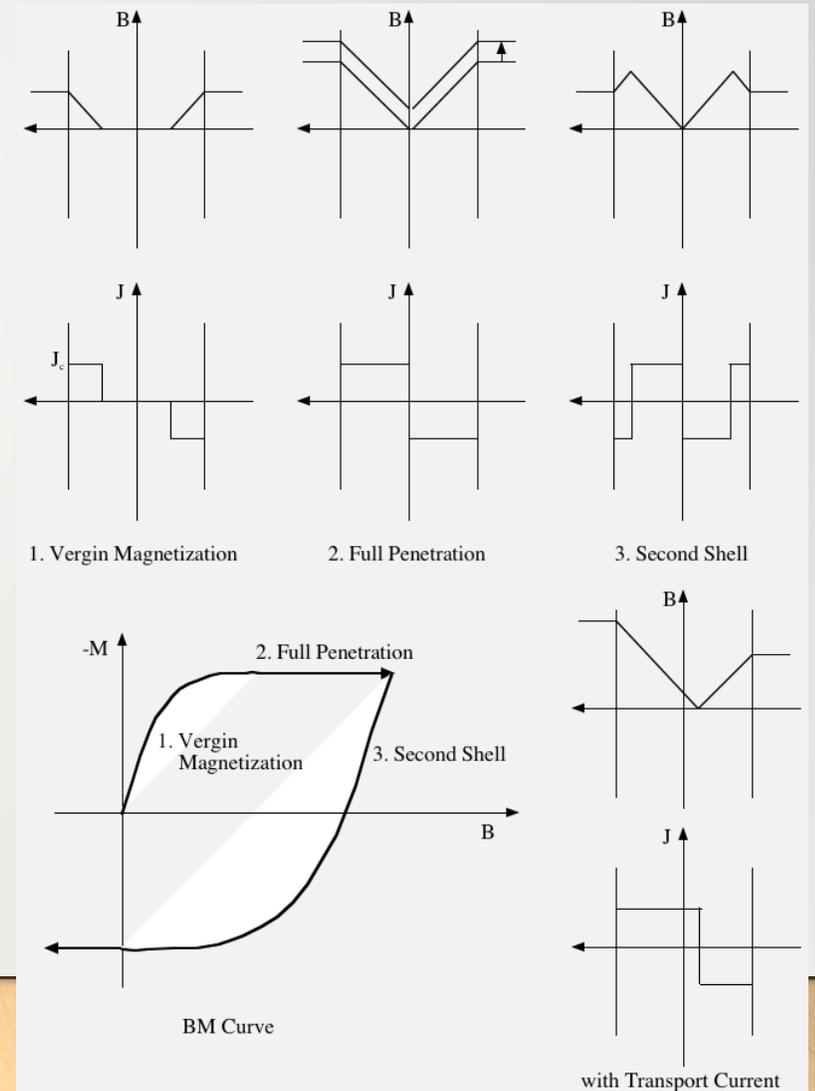
34

# Critical State Model: Magnetization

- By Bean and London
  - Flat panel model
  - thickness:  $a$
  - Magnetization:  $M$

$$-M = J_c a / 2$$

- proportional to  $a$  and  $J_c$
- fine filament: smaller magnetization



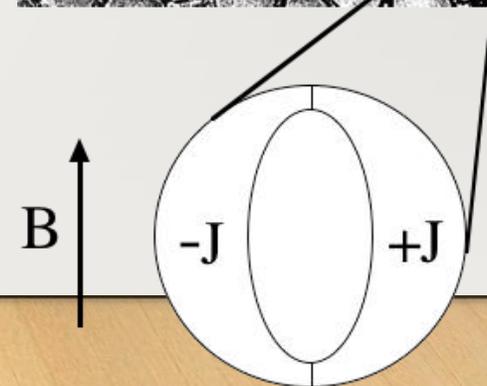
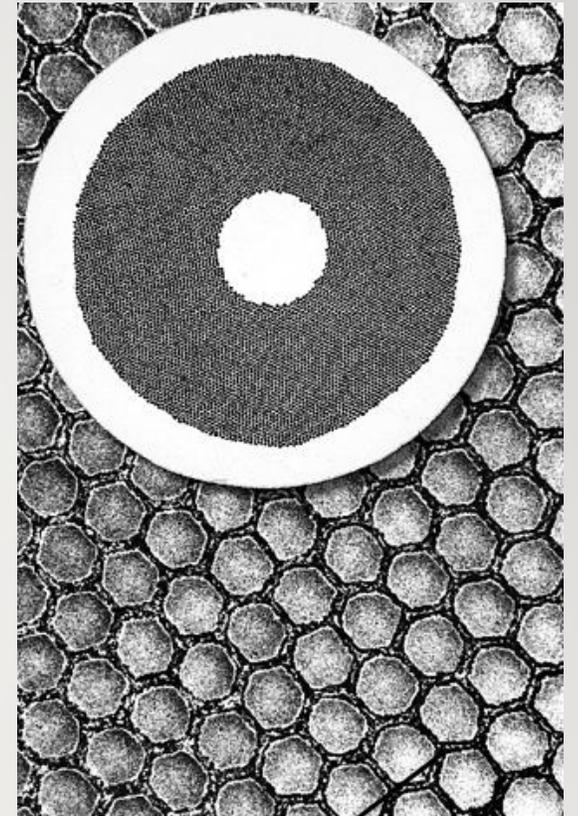
# 35 | Multi-filament composite

- Wire magnetization
  - Filament magnetization

$$M = \begin{cases} \frac{4}{3p} m_0 J_c r_f \left[ 1 - \left( 1 - \frac{B_0}{B_p} \right)^3 \right] & (B_0 < B_p) \\ \frac{4}{3p} m_0 J_c r_f & (B_0 \geq B_p) \end{cases}$$

where  $B_p = 2m_0 J_c r_f / p$

- Magnetization proportional to filament size



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# AC-Loss

- Hysteresis loss
  - Roughly proportional to saturated magnetization:  $M_s$   
( $\propto \mu_0 a J_c$ )
  - Proportional to field change
- Coupling loss
  - Roughly proportional to time constant:  $\tau$ 
    - For strand  $\tau \propto L_p^2 / \rho$  ( $L_p$ : twist pitch)
  - Proportional to field change square

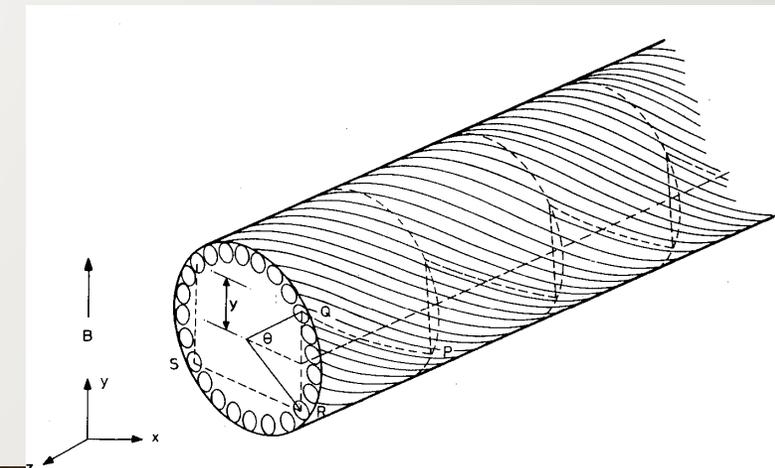
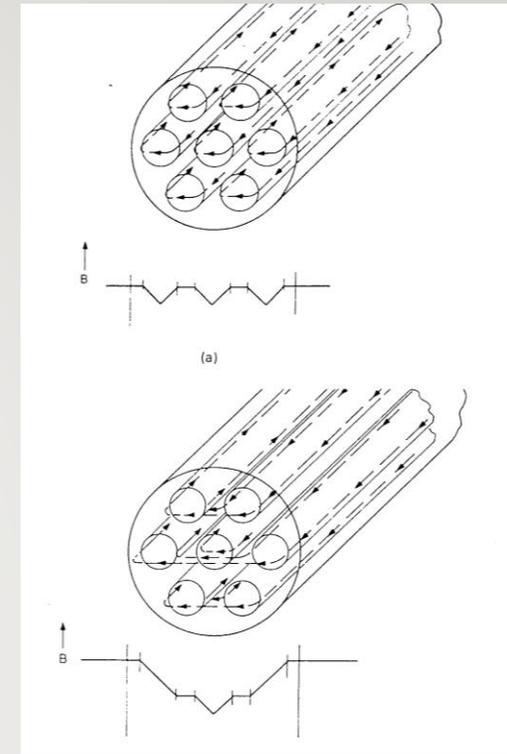
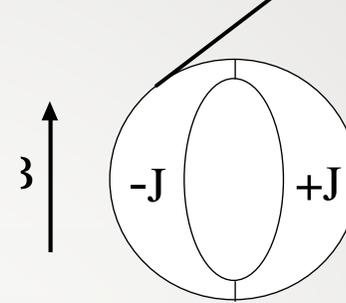
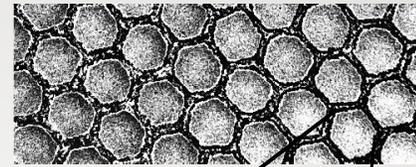
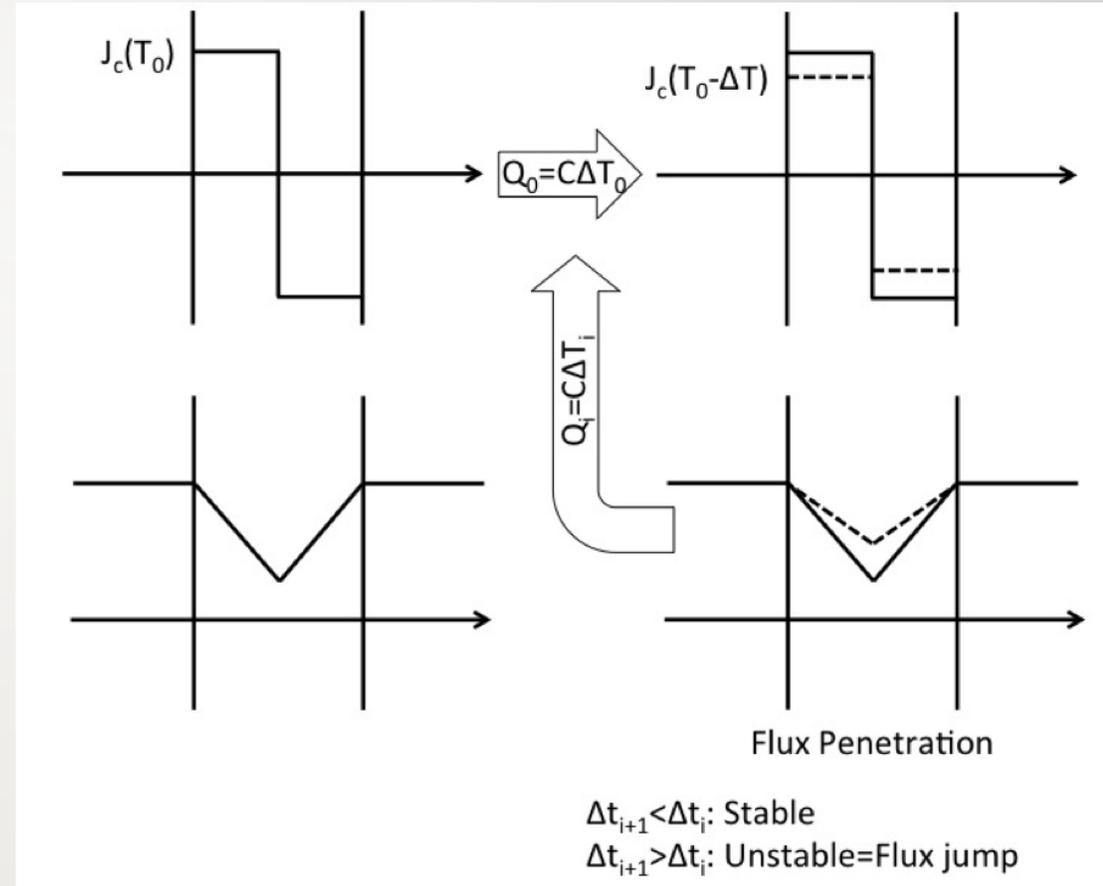


Fig. 8.15. Twisted filamentary composite in changing transverse field showing zig-zag path used to calculate flux linkages.

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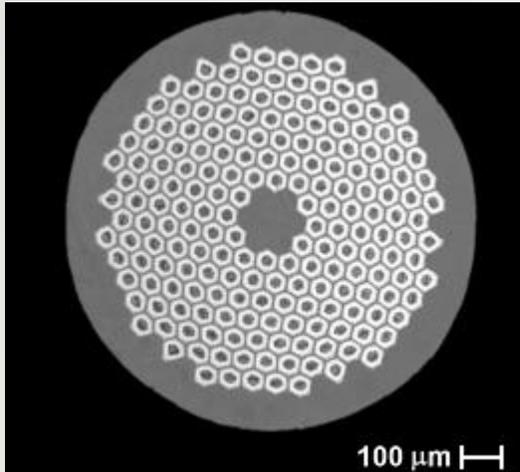
# Flux Jump

- Magneto-thermal instability
- Observed with
  - Large filament
  - High  $J_c$
  - Bad RRR
  - Low temperature operation

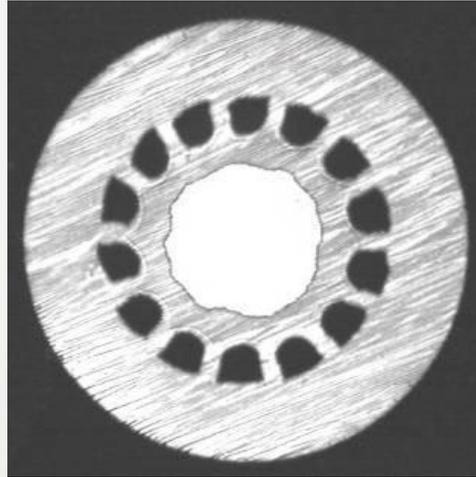


# Advanced Conductors

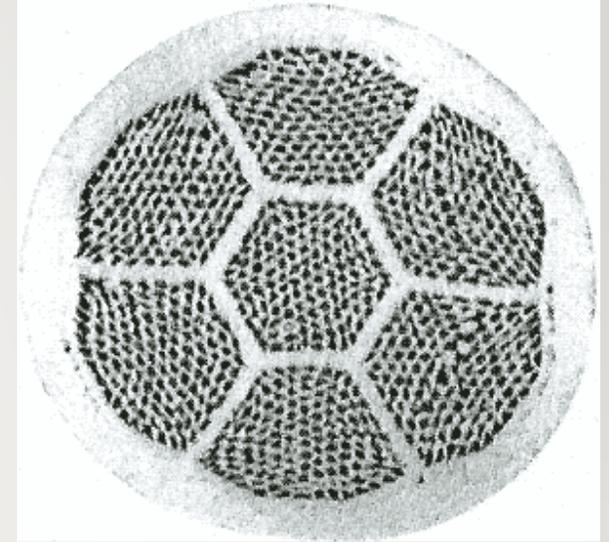
$\text{Nb}_3\text{Sn}$



$\text{MgB}_2$

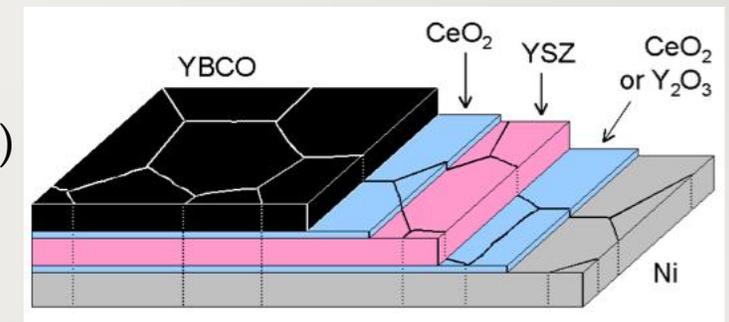


Bi-2212



Coated conductor. (IBAD, RABiTS or ISD)

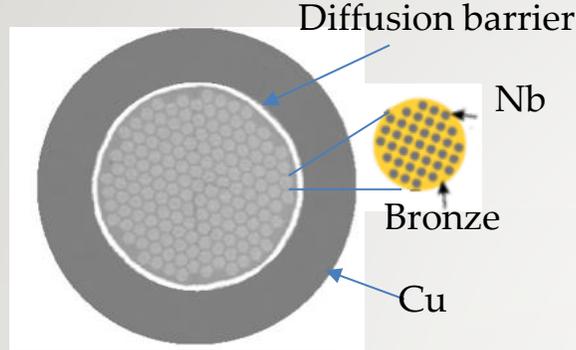
Bi-2223



# Nb<sub>3</sub>Sn industrial wires

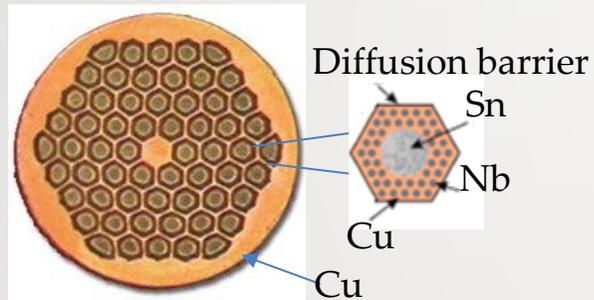
39

## ➤ Bronze Route



Small filaments ( $\Phi < 5 \mu\text{m}$ )  
Jc limited ( $\sim 1\text{-}1.2 \text{ kA/mm}^2$  @ 12 T, 4.2 K) by the solubility of Sn in Cu ( $\sim 15.5 \text{ wt } \%$ )

## ➤ Internal Tin

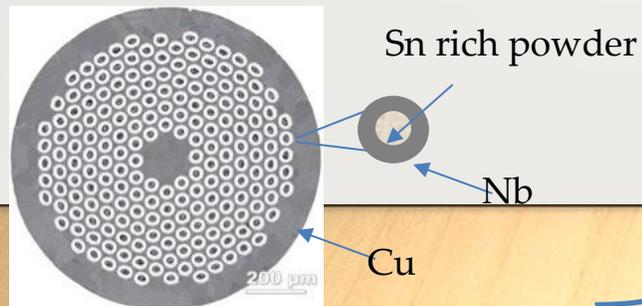


RRP®

*The Business of Science®*

High Jc Nb<sub>3</sub>Sn wires for accelerator magnets

## ➤ Powder In Tube



PIT

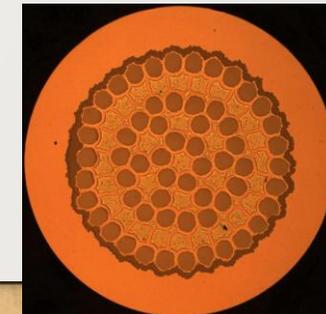
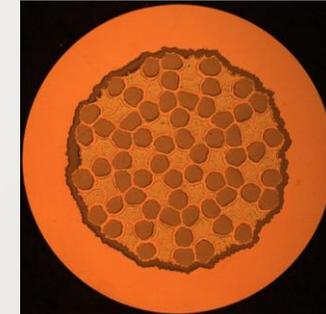
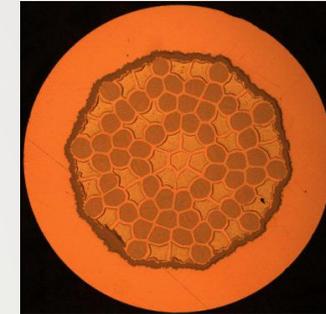
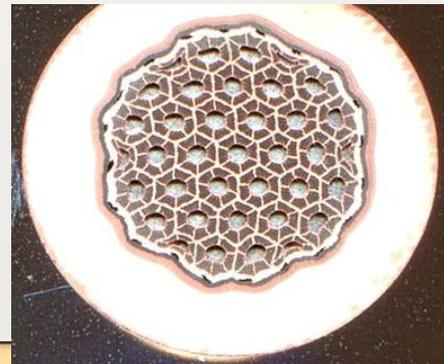
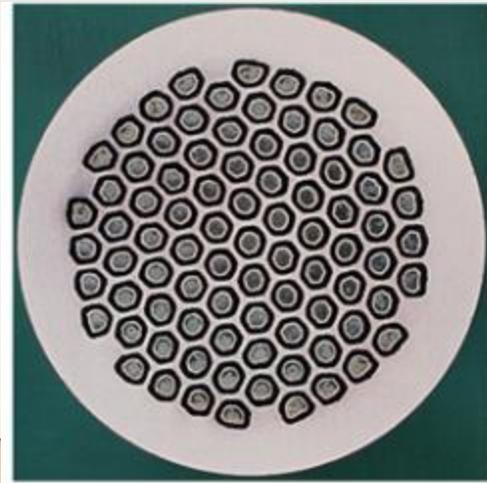
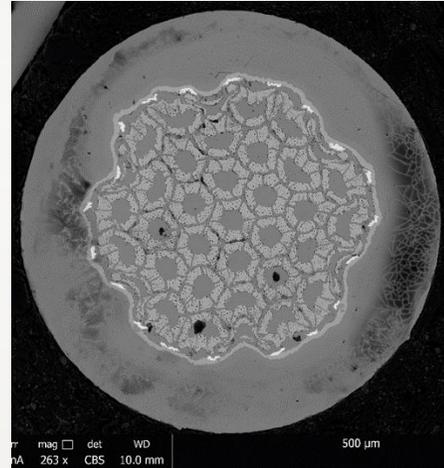
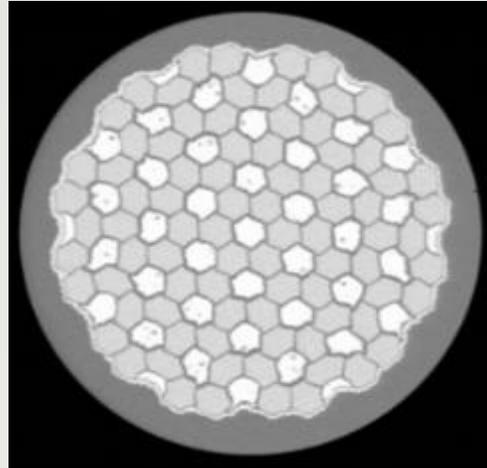
40

# Nb<sub>3</sub>Sn development CERN world wide collaboration

**KOBELCO**  
KOBEL STEEL GROUP

**S JASTEC**  
SUPERCONDUCTOR

**FURUKAWA  
ELECTRIC**



NATIONAL NUCLEAR CORPORATION "ROSATOM"



FUEL COMPANY OF ROSATOM  
**TVEL**

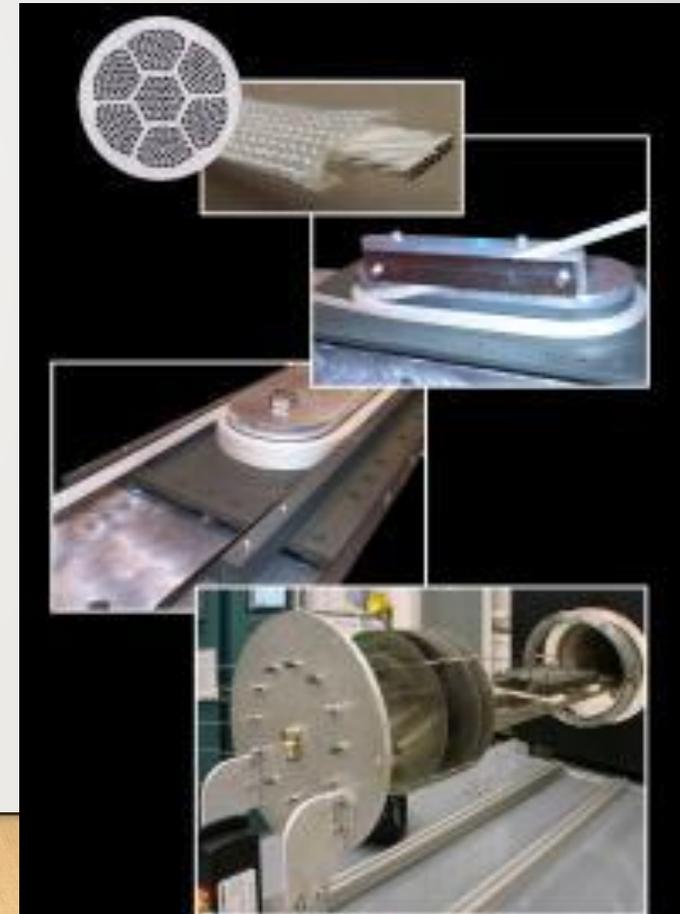


**BOHVAR INSTITUTE OF  
INORGANIC MATERIALS  
JSC VNIINM**

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# HTS conductor BISCO

- Bi2223
  - Tape Shape
  - Magnetic Anisotropy
  - React and Wind
  - Mechanical Reinforcement
- Bi2212
  - Developed by Japanese Company: improved by USA
  - Round Wire
    - Rutherford Cable
  - Wind and React
    - under high pressure O<sub>2</sub>



# HTS Conductor YBCO Coated Conductor

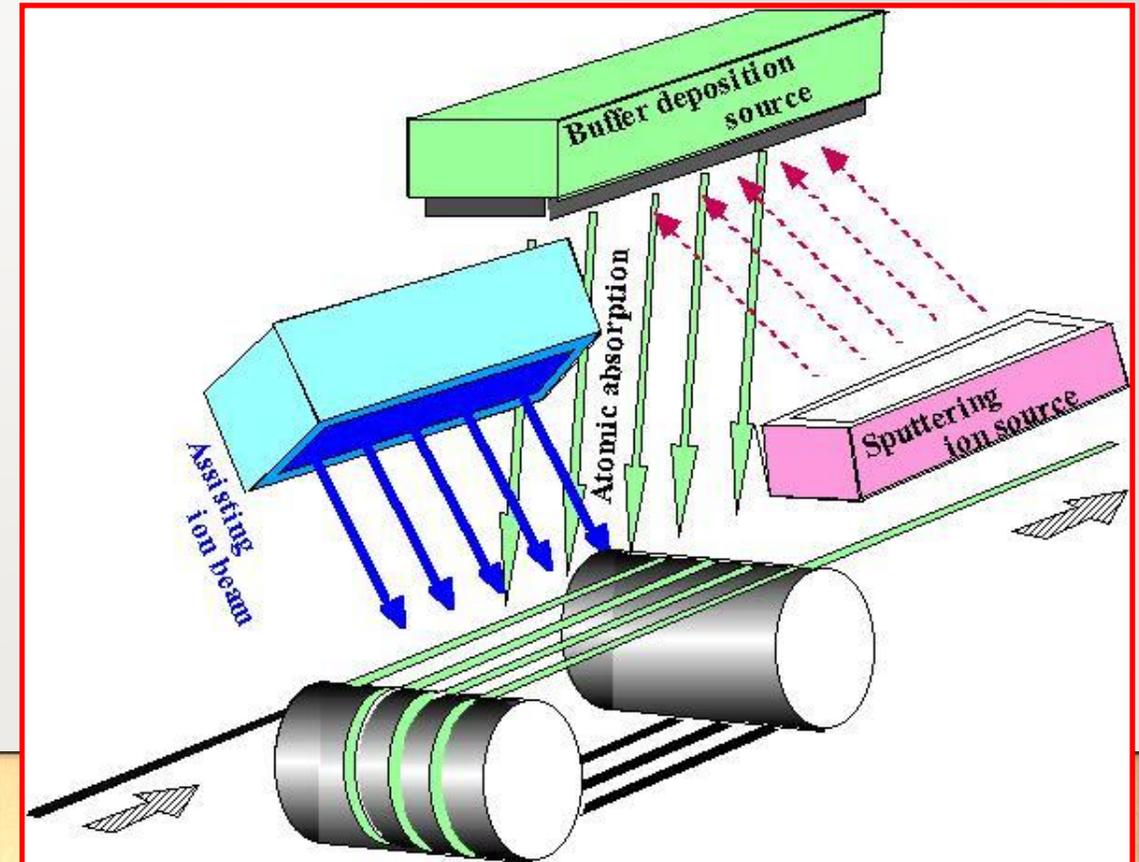
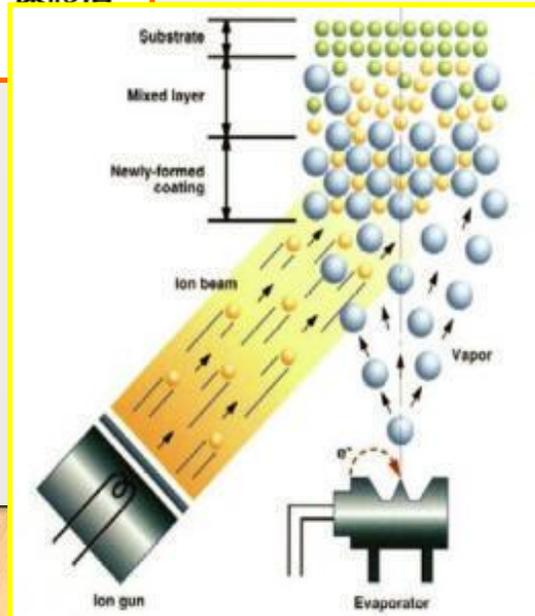
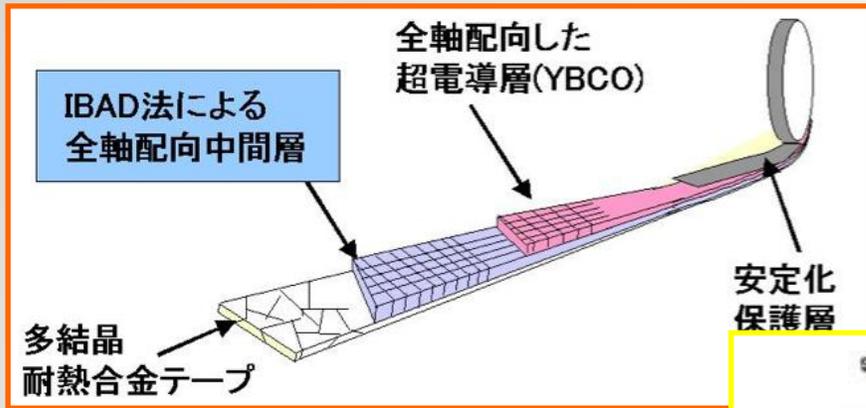
Tape Shape: Magnetic Anisotropy

Wide superconductor: Big Shielding current against perpendicular field

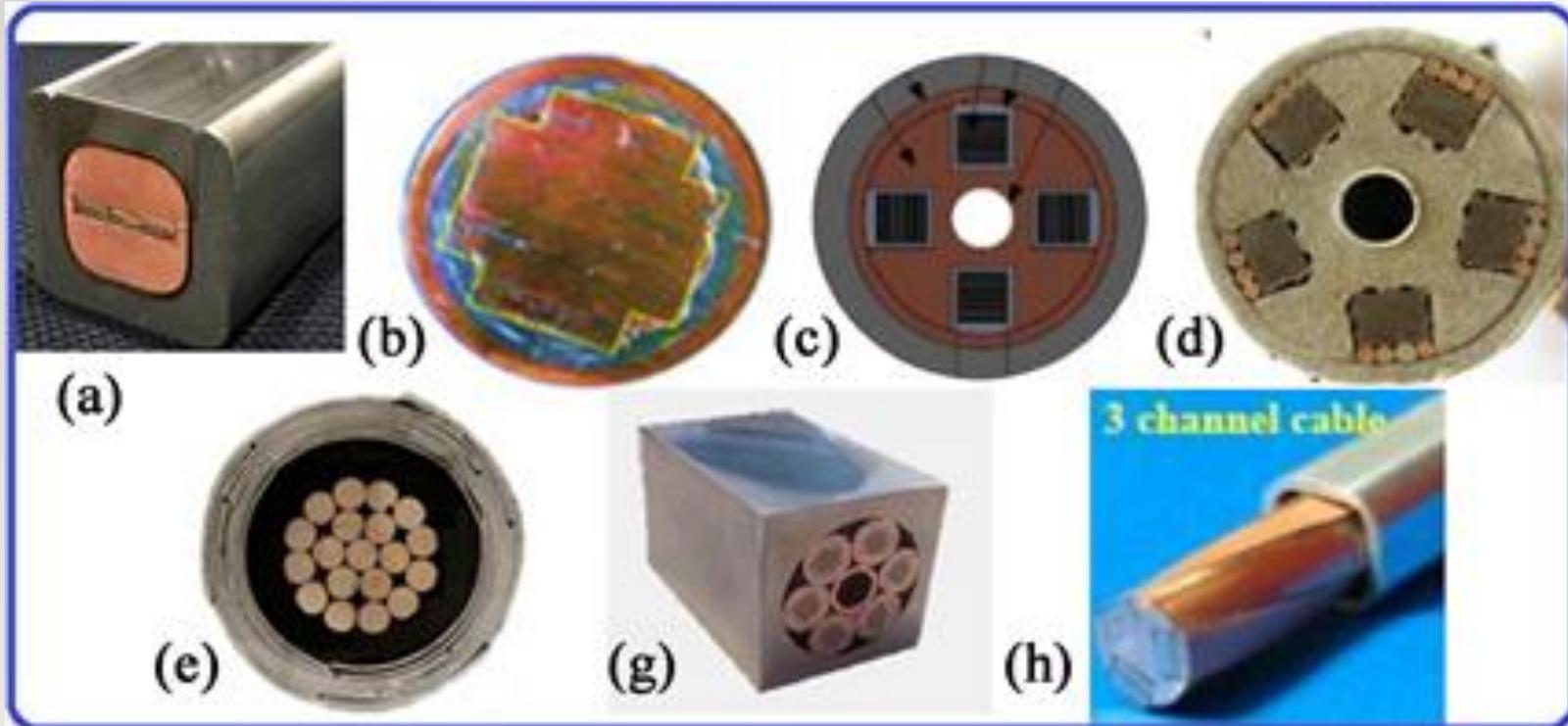
React and Wind

Strong against longitudinal force

Weak against tearing force



# HTS Cables



Cross sections of various HTS high-current Cables

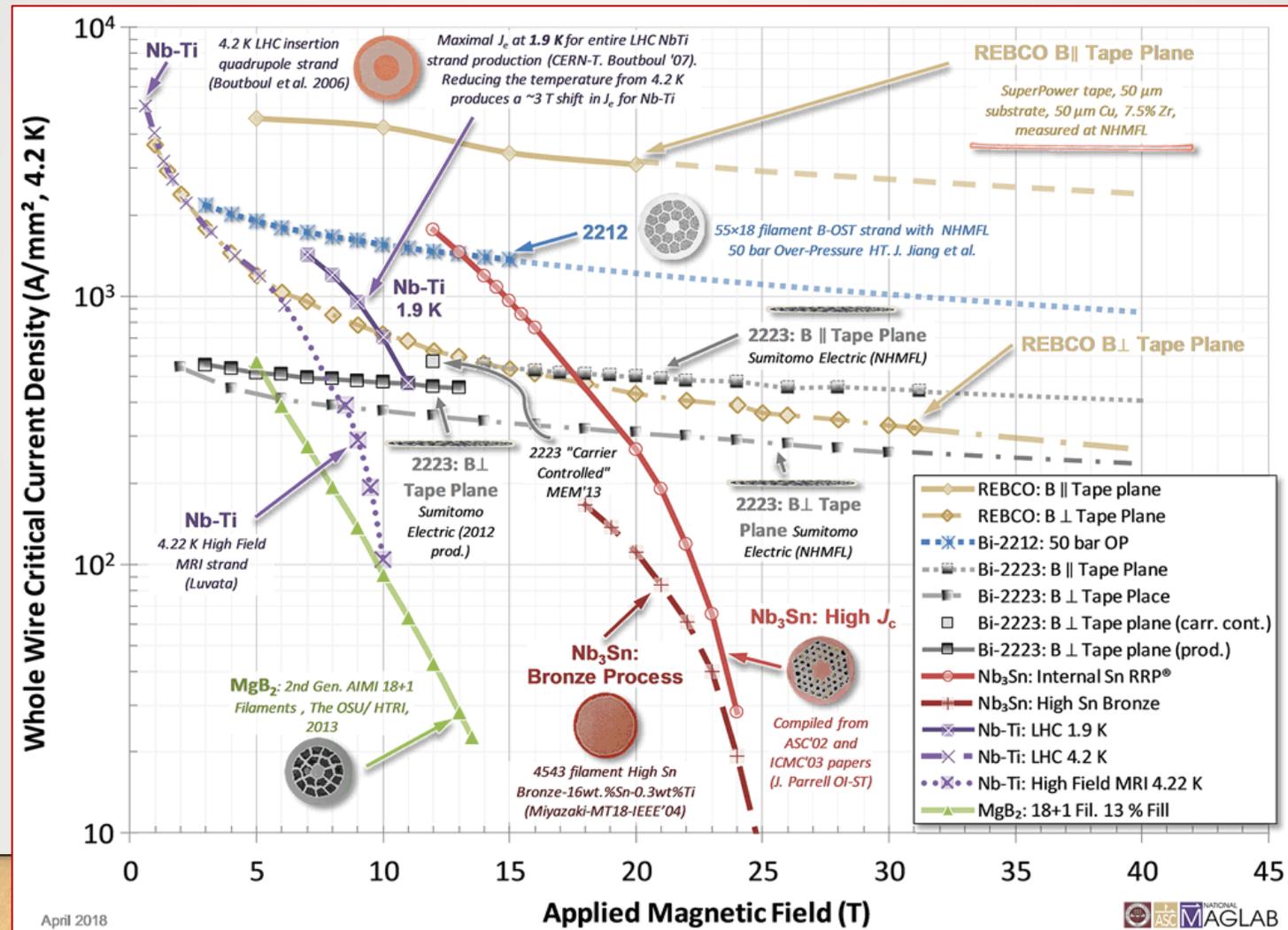


Roebel Cable



CORC Cable

# Performance of Superconductors



# Item

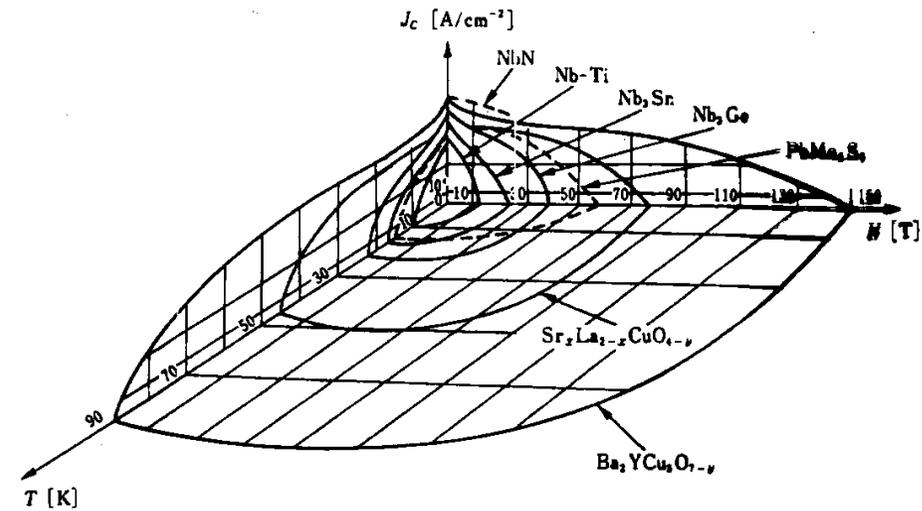
- Superconductivity
- Superconducting Wire and Cables
- Stability and Quench Protection
- Summary

# Quench ?

- Superconductor
  - Critical Surface
  - Exceeding critical surface = Quench
- Real operation point must have margin
  - Temperature margin  $\Leftrightarrow$  Heat disturbance
  - Disturbance : Time and Space distribution

	Point	Distributed
Transient	Joules	Joules/m <sup>3</sup>
Continuous	Watts	Watts/m <sup>3</sup>

- Cooling Condition : Direct by coolant or indirect
- Quench  $\leftarrow$  Depends on those conditions
  - **Quench Stability: important research item**



典型的な超伝導体の臨界特性<sup>[2]</sup>

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# Adiabatic VS Distributed Disturbance

- Transient
    - Heat capacity
- $$Q \Leftrightarrow C\Delta T$$
- Heat capacity at low temperature is very low
- Continuous
    - For adiabatic condition it raise temperature continuously

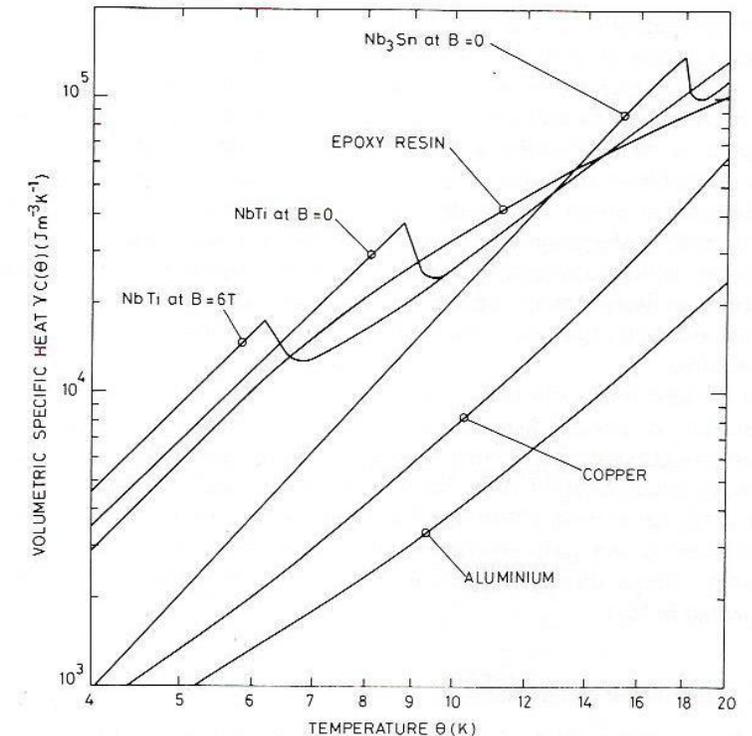


Fig. 5.3. Volumetric specific heat of some materials commonly used in magnets. Data for NbTi are average values based on measurements by J. Corsan, National Physical Laboratory (private communication) and R. Lagnier and E. Bonjour, CEN Grenoble (private communication). Data for  $\text{Nb}_3\text{Sn}$  is from Grigsby D. L. (US Department of Commerce NTIS Report AD-838469) and remaining data is from Johnson (1961).

# Adiabatic vs Point Disturbance

Heat conduction

vs

Joule heating

$$2kA(q_c - q_0)/l = J_c r Al$$

Minimum Propagating Zone (MPZ)

$$2kA(q_c - q_0)/l = I^2 J_m r [J_m - J_c(q)] / (1 - l) Al$$

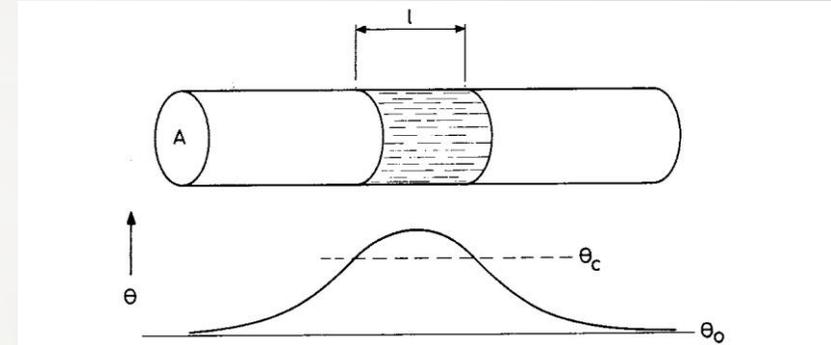
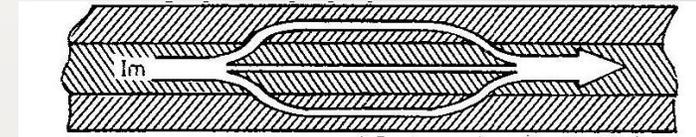


Fig. 5.8. A small normal zone created by a point disturbance in a current-carrying wire.



Current Sharing in Composite Conductor

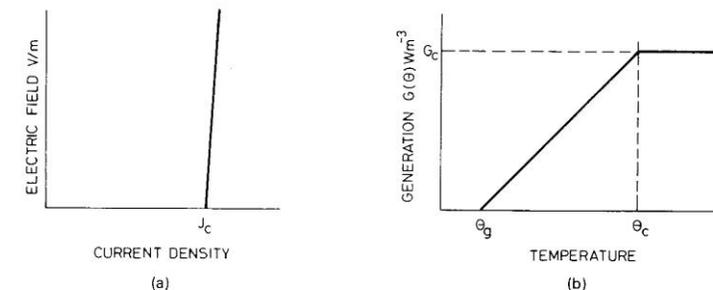


Fig. 5.9. (a) The development of flux flow resistance in a superconductor at current densities above critical; (b) temperature dependence of power generation in a composite conductor.

# With Cooling

- Disturbance → Continuous or Transient?
- Continuous
  - Cooling power vs Heat input
- Transient
  - Heat capacity vs Heat Disturbance → Size of normal zone → Joule heating
  - Cooling power vs Joule heating (MPZ)

# Cooling Condition: Pool Boiling Helium

- Heat Transfer
  - Nucleus boiling
    - Wet surface
    - Critical heat transfer
    - Sub-cool  $\rightarrow$  some margin to start boiling
  - Film boiling
    - Dry surface
    - Transition from nucleus boiling
    - Significantly lower heat transfer

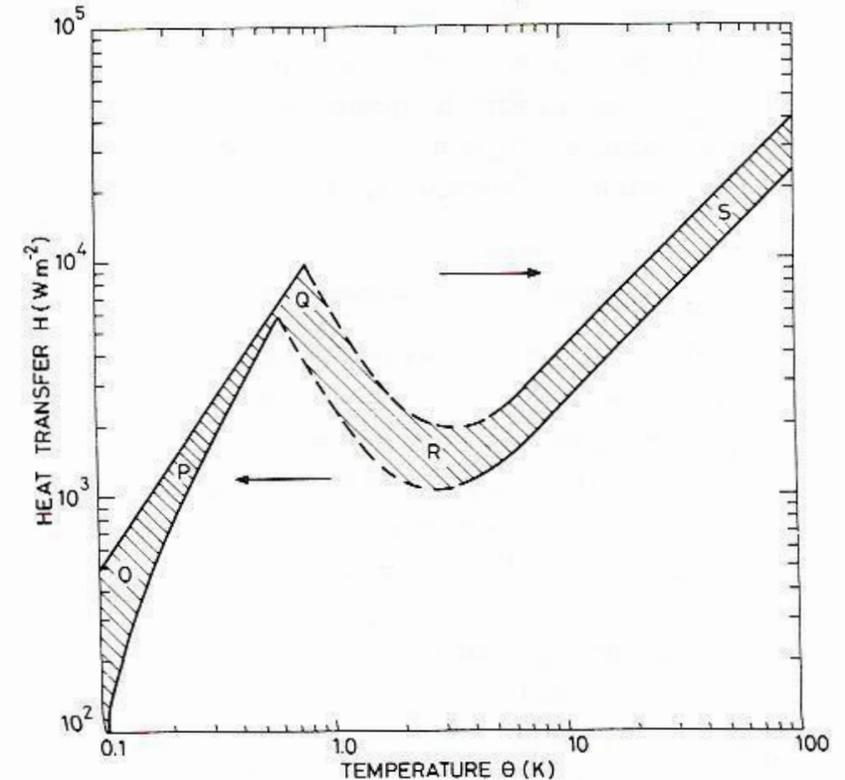
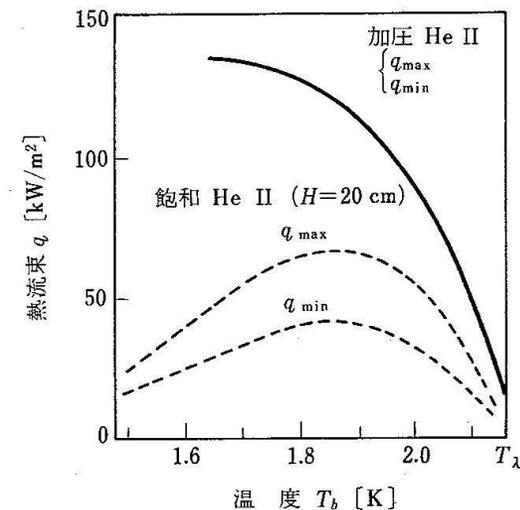


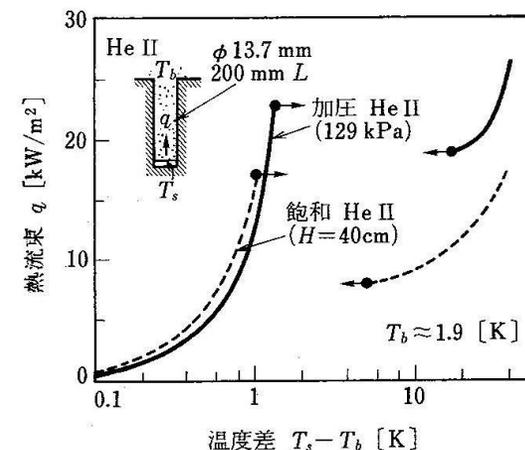
Fig. 6.3. Heat transfer from a metal surface to liquid helium boiling at 4.2 K under normal atmospheric pressure (shaded areas indicate spread of results).

# Cooling Condition: Super Fluid Cooling

- Heat Transfer
  - Kapitza resistance
    - Material boundary → Heat resistance
    - By Phonon Scattering
  - Heat Transfer in Helium
    - Very big
    - No nucleus boiling
      - Kapitza resistance dominate heat transfer upto film boiling
- Saturated vs Pressurized
  - Pressurized Super Fluid
    - Bigger critical heat transfer
    - The best condition is
      - 1.8~1.9K
      - 0.1 MPa



■ 5・54  $q_{max}$  および  $q_{min}$  の温度依存性 [Kobayashi, et al. 1980]



■ 5・55 加圧 He II の熱伝達特性 [Van Sciver 1978]

# Cooling Condition: Supercritical

- Forced Flow Cooling
- Supercritical
  - No boundary between liquid and gas
  - No phase transition
- Heat Capacity
  - Peak at 0.3 MPa, 5~6K
- Heat Transfer
  - Peak around 0.3 MPa, 5~6K
- ~0.3MPa Operation
  - Maximum temperature 5~6K

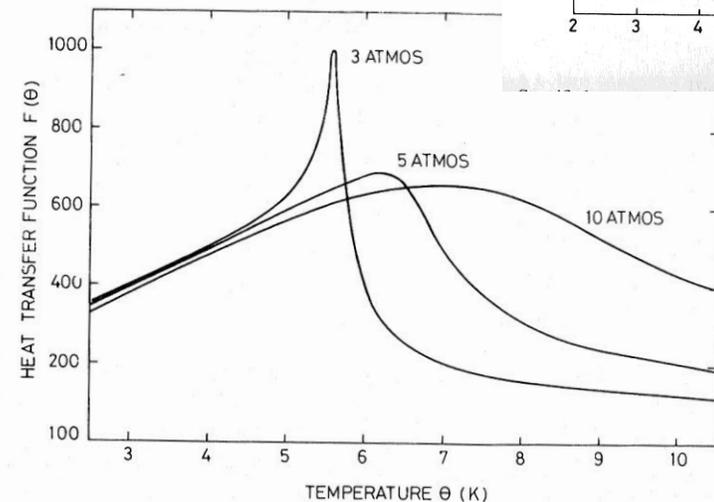
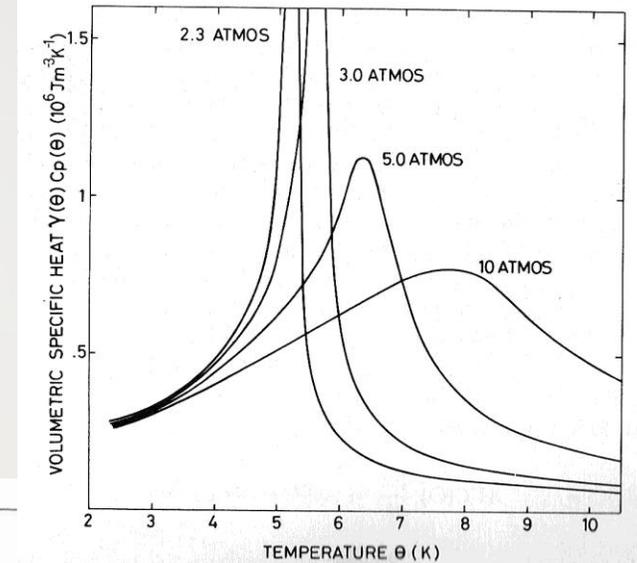


Fig. 6.25. Heat transfer function defined by eqn (6.52) (data from McCarty 1972).

# Cooling Condition vs Distributed Disturbance

- Perfect stabilization : Stekley

$$Ph(q - q_0) = I^2 J_m r [J_m - J_c(q)] / (1 - I) A$$

$$\text{if } J_c(q) = J_{c0} (q_c - q) / (q - q_0)$$

$$1 = \frac{I^2 J_c^2 r A}{(1 - I) Ph(q_c - q_0)} = \frac{G_c A}{Ph(q_c - q_0)} = a : \text{Stekly parameter}$$

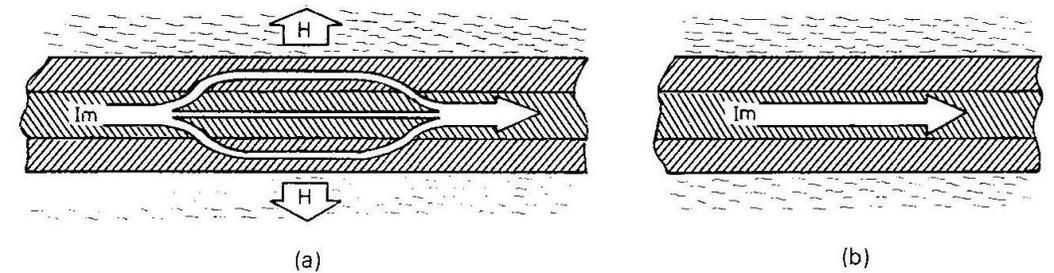


Fig. 6.1. (a) Cryogenic stabilization; following a disturbance current shares between copper and superconductor; ohmic heat generation is transferred to the coolant; (b) if the available cooling exceeds generation, temperature falls and current returns to the superconductor.

# Cooling Condition vs Semi-Distributed Disturbance

- Madock theorem
  - Cooling along conductor length
  - Basic equation

$$\frac{d}{dz} \left( Ak(q) \frac{d}{dz} \right) = PH(q) - AG(q)$$

- Equal-Area Theorem
  - Heat transfer curve vs Joule heating curve

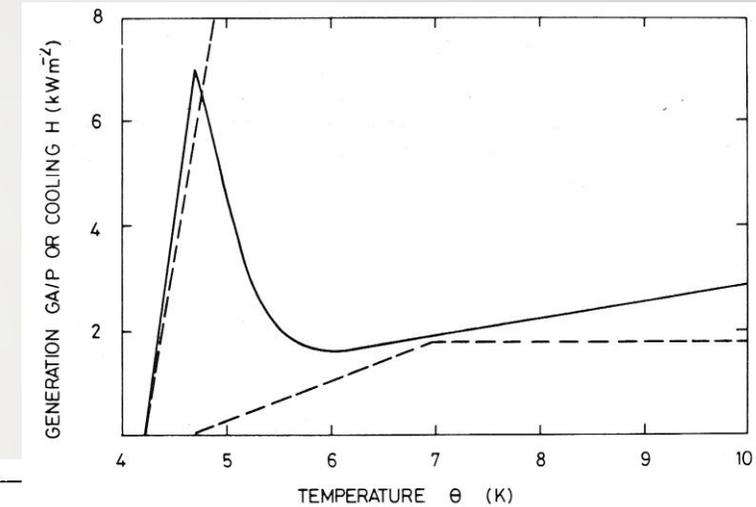
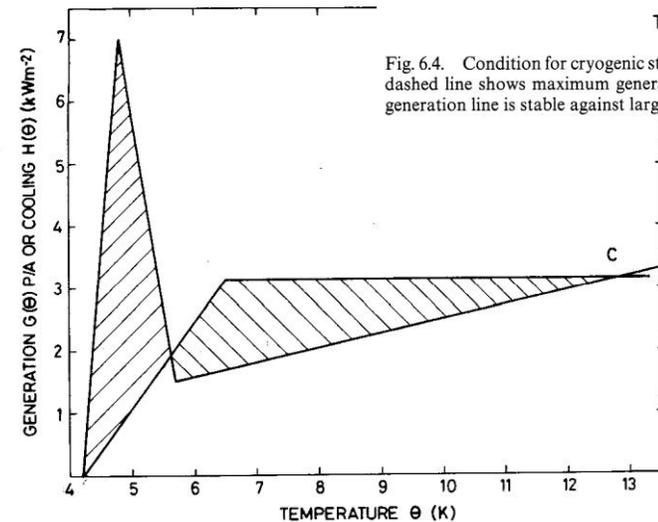


Fig. 6.4. Condition for cryogenic stability with boiling heat transfer (no heat conduction): upper dashed line shows maximum generation for stability against small disturbances; lower dashed generation line is stable against large disturbances.



6.6. Equal-area condition for cryogenic stability of a normal zone with cold superconducting regions at either end (drawn for NbTi at 6 T, i.e.  $\theta_c = 6.5$  K).

# Cooling Condition vs Point Disturbance

- Extend Madock Theorem
  - Small normal zone
  - Shrink of normal zone
  - MPZ
  
- For Real magnets
  - 3D anisotropy
    - Along conductor
    - Turn to turn
    - Layer to layer
  - Basic equation

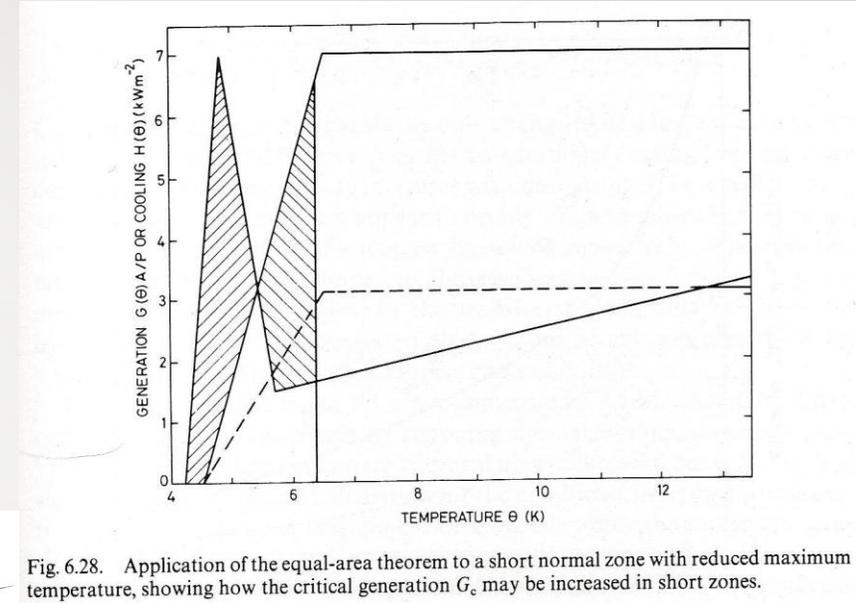
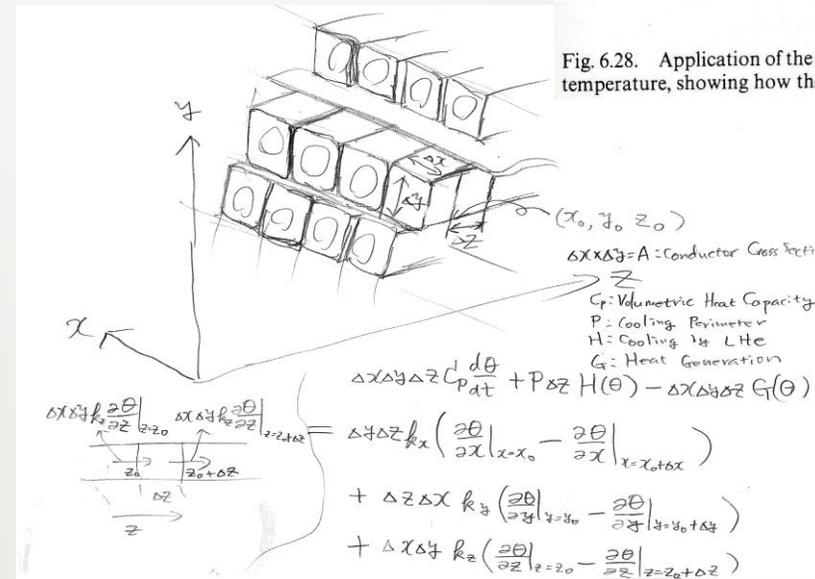


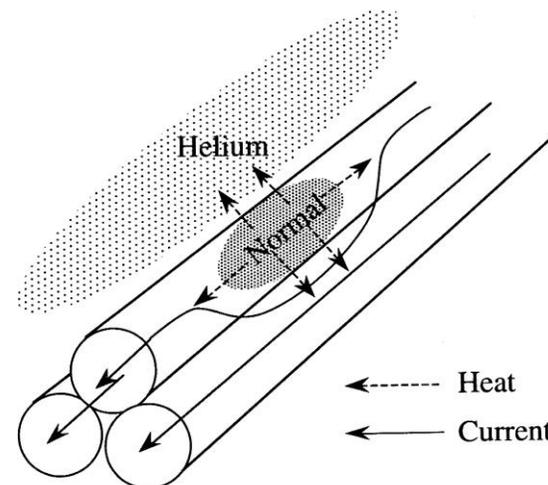
Fig. 6.28. Application of the equal-area theorem to a short normal zone with reduced maximum temperature, showing how the critical generation  $G_c$  may be increased in short zones.



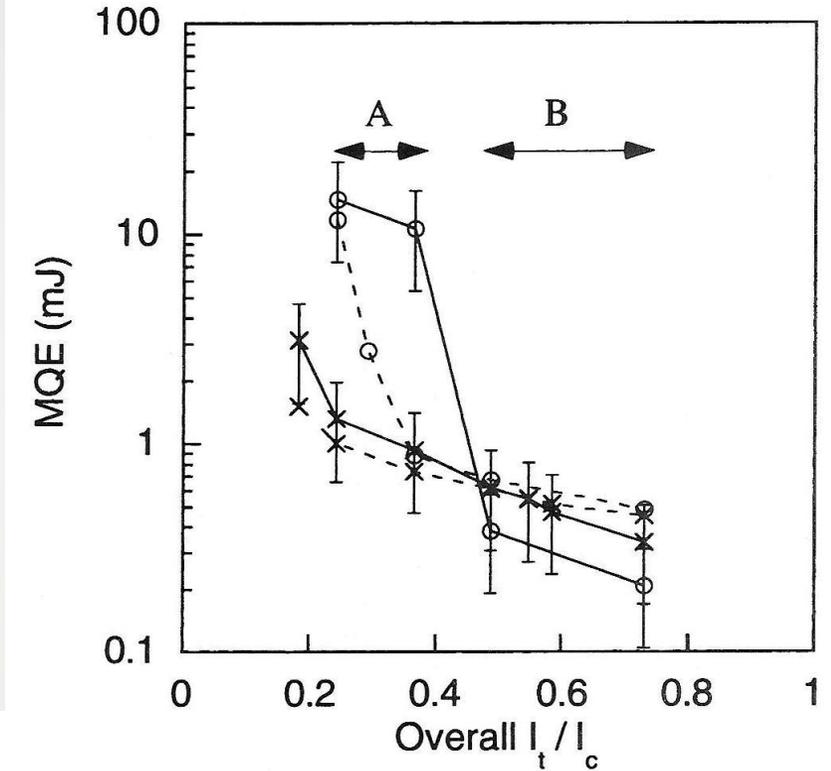
$$\frac{\partial}{\partial x} \left( k_x(q) \frac{\partial}{\partial x} \right) + \frac{\partial}{\partial y} \left( k_y(q) \frac{\partial}{\partial y} \right) + \frac{\partial}{\partial z} \left( k_z(q) \frac{\partial}{\partial z} \right) = C_p \frac{\partial q}{\partial t} + \frac{P}{A} H(q) - G(q)$$

# Stability on Cable

- Cable = Multiple Strands
  - Inter-strand heat transfer
  - Inter-strand current sharing
- Better stability
  - with lower inter-strand resistance
  - AC-loss will increase



**Figure 1** Current transfer and thermal conduction between strands following formation of local normal zone



**Figure 9** MQE vs.  $I_t / I_c$ . Lines with open circles denote MQE of triplex cable against local disturbance, and lines with crosses denote MQE of single strand. Solid lines denote experimental results, and broken lines denote numerical results where contact electrical conductance and contact thermal conductance are  $1.36 \times 10^5$  S/m and 2 W/(Km), respectively. In region A, current re-distribution is observed in recovery process, and in region B, it is not observed

# Quench and Protection

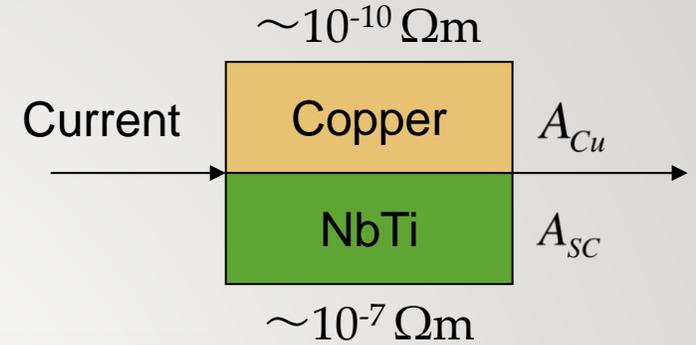
- Quench = Abrupt phase transition to normal
- Current flow in Copper → Still fast temperature rise
- Temperature rise depends: MIITs

$$\int_{t_{quench}}^{t_{end}} I^2 dt$$

$$= \int_{T_0}^{T_{max}} \frac{C_{pCu} A_{Cu} + C_{pSC} A_{SC}}{r_{Cu} / A_{Cu}} dT$$

- MIITs/I<sup>2</sup>: time to reach a temperature

Cu/NbTi=1, 400 A/mm<sup>2</sup> →

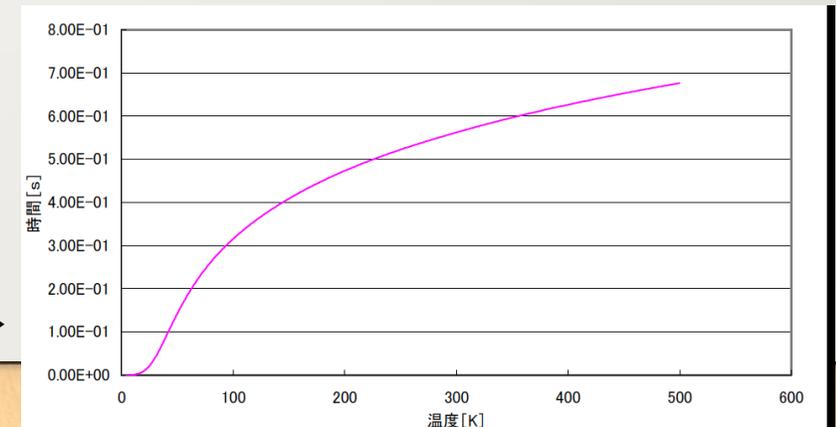


$$R = r_{Cu} l / A_{Cu}$$

$$G = I^2 r_{Cu} l / A_{Cu}$$

$$= \Delta H$$

$$= (C_{p-Cu} A_{Cu} + C_{p-SC} A_{SC}) l \Delta T$$



Rapid detection of quench and shutting down current is needed

# Quench Propagation

- Quench Propagation
  - Basic Equation (same as stability)

$$\frac{\partial}{\partial x} \left( k_x(q) \frac{\partial}{\partial x} \right) + \frac{\partial}{\partial y} \left( k_y(q) \frac{\partial}{\partial y} \right) + \frac{\partial}{\partial z} \left( k_z(q) \frac{\partial}{\partial z} \right) = C_p \frac{\partial q}{\partial t} + \frac{P}{A} H(q) - G(q)$$

- One dimensional adiabatic assumption

$$v_{ad} = \frac{J}{gC} \left\{ \frac{rk}{q_s - q_0} \right\}^{\frac{1}{2}}$$

$$= \frac{J}{gC} \left\{ \frac{L_0 q_s}{q_s - q_0} \right\}^{\frac{1}{2}}$$

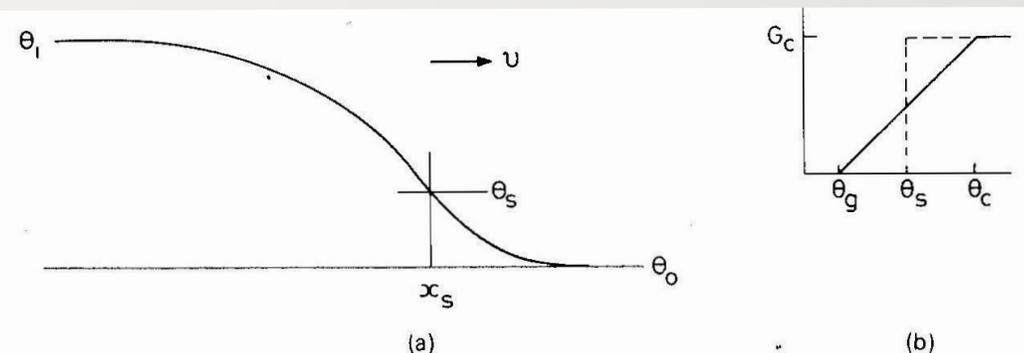


Fig. 9.3.(a) Temperature profile at the boundary of a normal zone advancing from left to right; (b) approximation of the generation curve by an abrupt transition.

## 59 | Quench Stability and Protection Summary

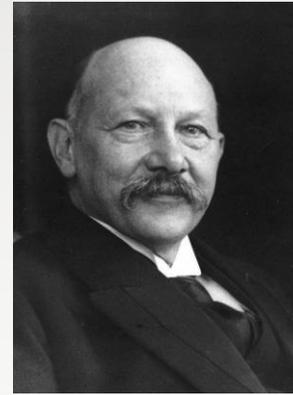
- Stability Analysis
  - Well established
  - Surface or boundary condition is not clear
- Steckley, Madock conditions
  - Give good stability margin
    - Lower current density
    - More aggressive condition was actually used
    - For HTS; large heat capacity give a huge stability
- Quench Protection
  - Quench detection and current shut down with in MIITs limit
  - For HTS due to low quench velocity quench detection is issue

# Item

- Superconductivity
- Superconducting Wire and Cables
- Stability and Quench Protection
- Summary

# 61 | Summary

- SC Conductor and Magnet
- Developed supplementary
- Reached to NbTi limit
- Higher Field by Nb<sub>3</sub>Sn (or IBS?)
  - FCC world wide development
- HTS (High Temperature Superconductor)
  - Still very expensive
  - Needs special usage
    - High Radiation
    - High Filed above 20 T
- **Game Change = New Chance!**



Kamerlingh Onnes  
Find Superconductivity



Alvin Tollestrup  
TEVATRON Magnet

Who's next  
YOU?

